

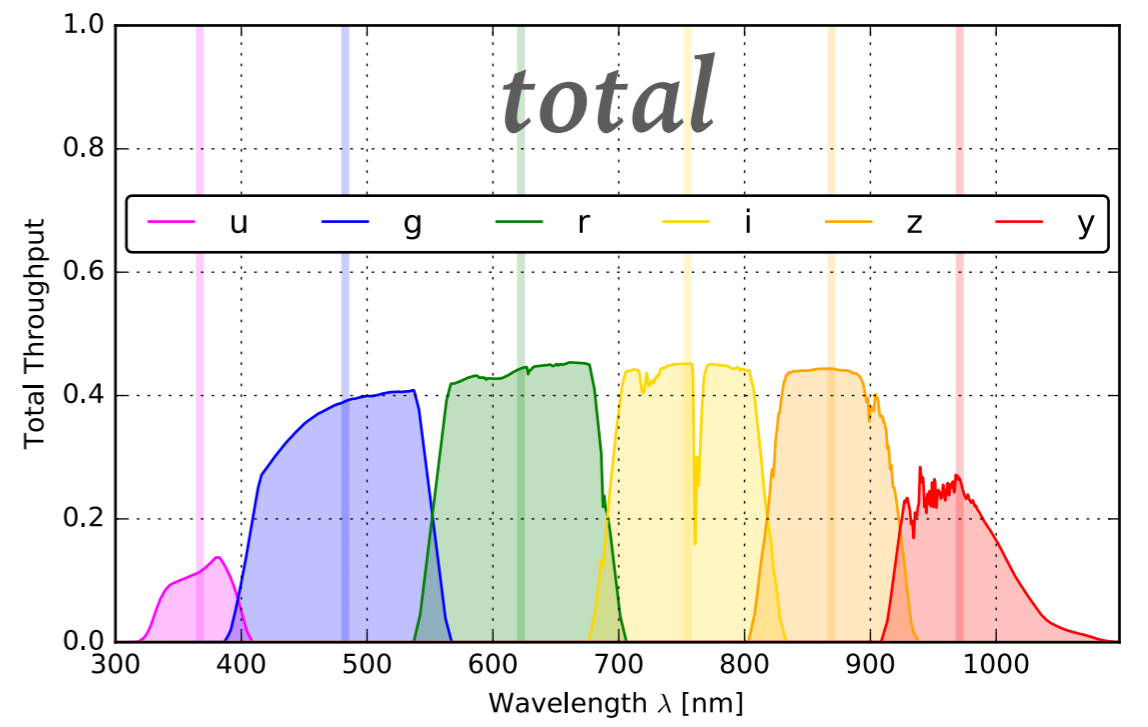
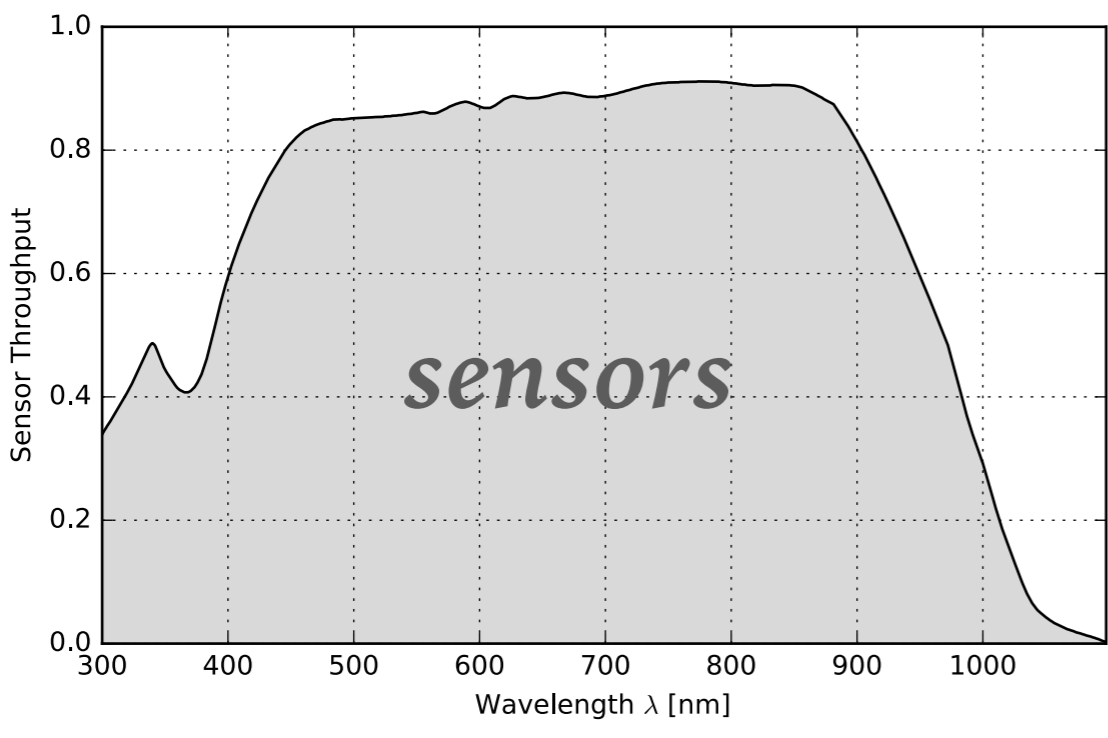
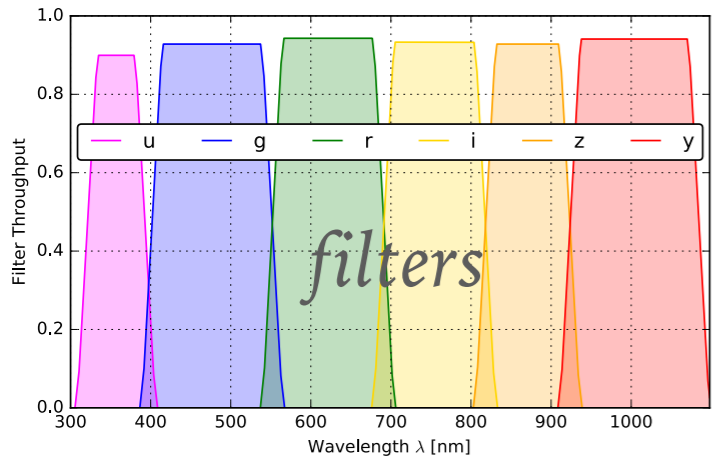
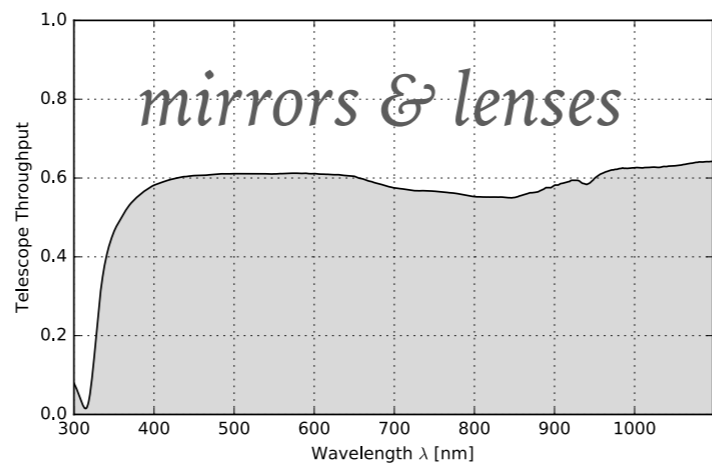
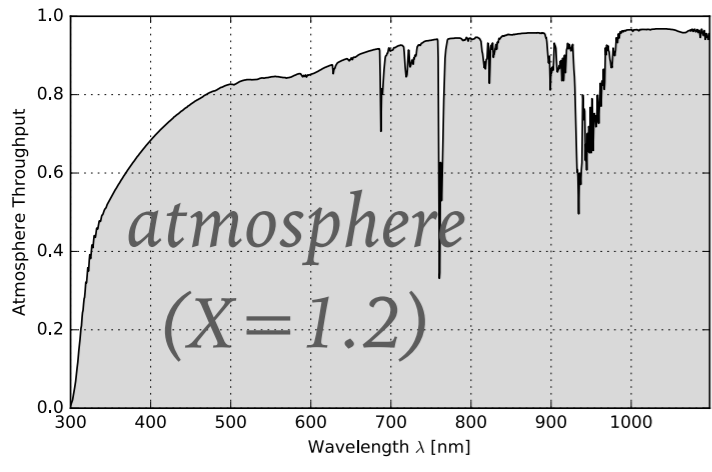
WAVELENGTH-DEPENDENT SENSOR RESPONSE

*Bela Abolfathi, David Kirkby,
Kristen McGee, Simona Murgia.*

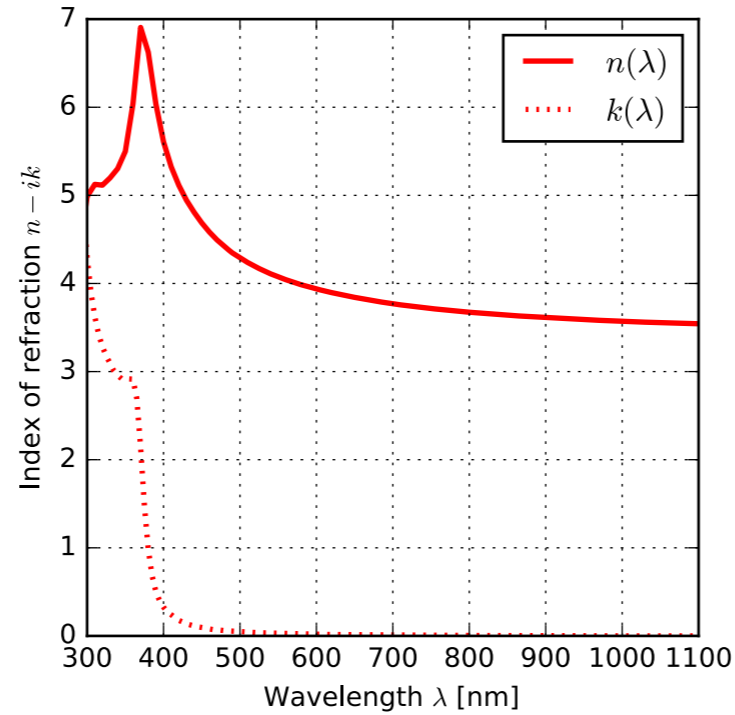
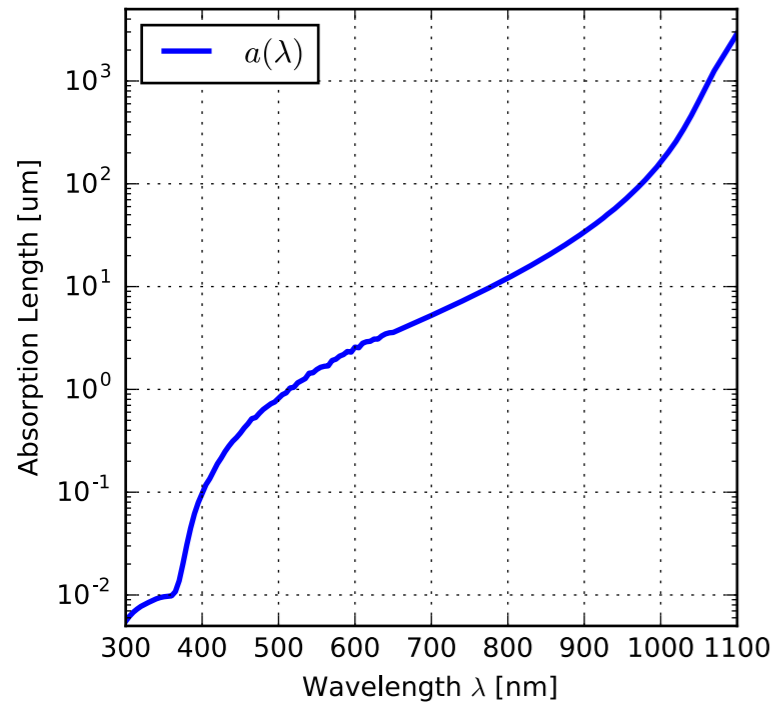
Univ. of California, Irvine

LSST DESC Collaboration Meeting, Oxford, July 2016.

LSST WAVELENGTH RESPONSE



SILICON OPTICAL PROPERTIES

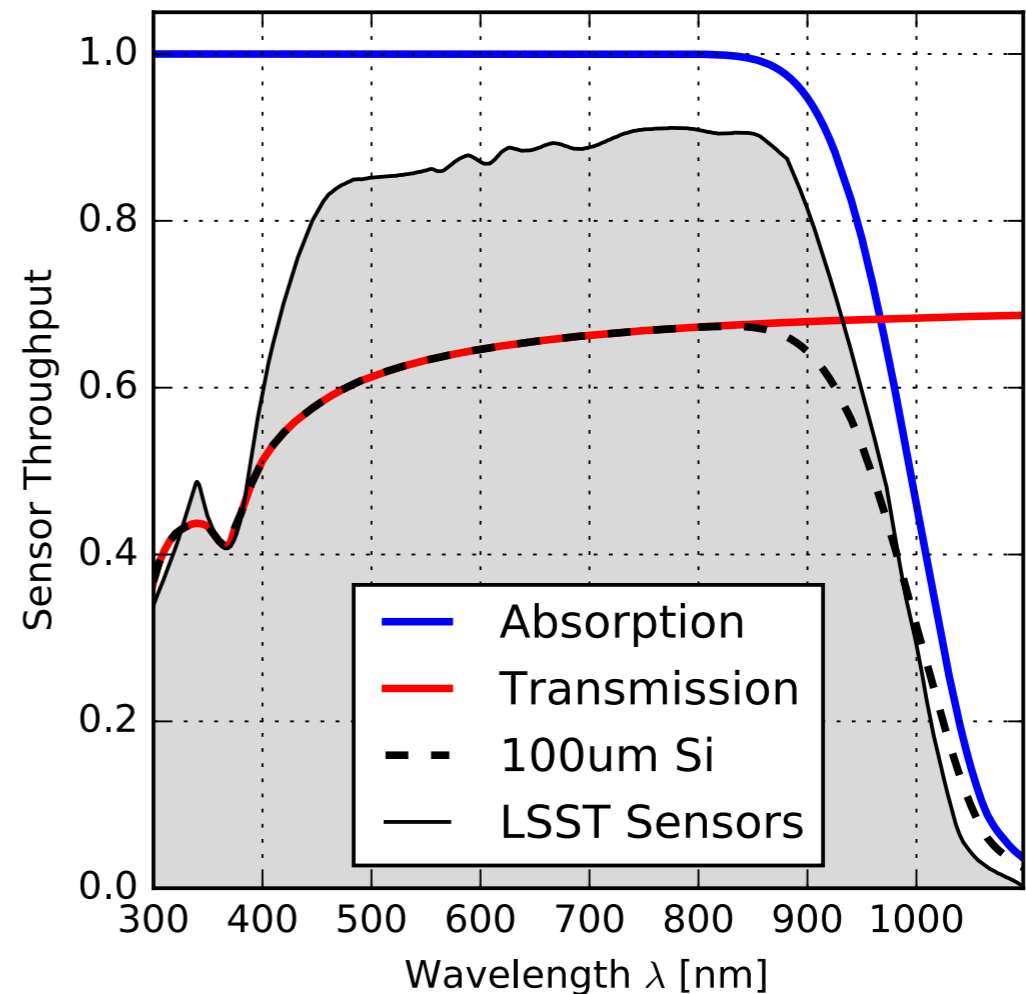


Absorption probability in $\Delta z = 100 \mu\text{m}$:

$$A(\lambda, \Delta z) = 1 - \exp(-\Delta z/a(\lambda))$$

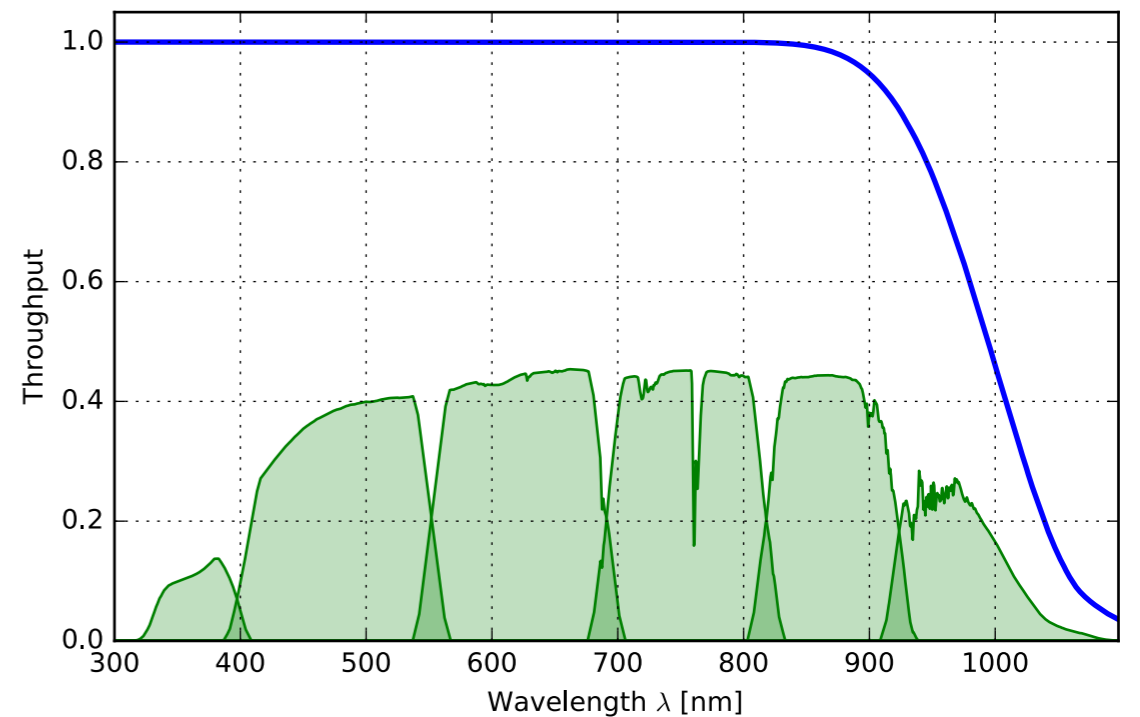
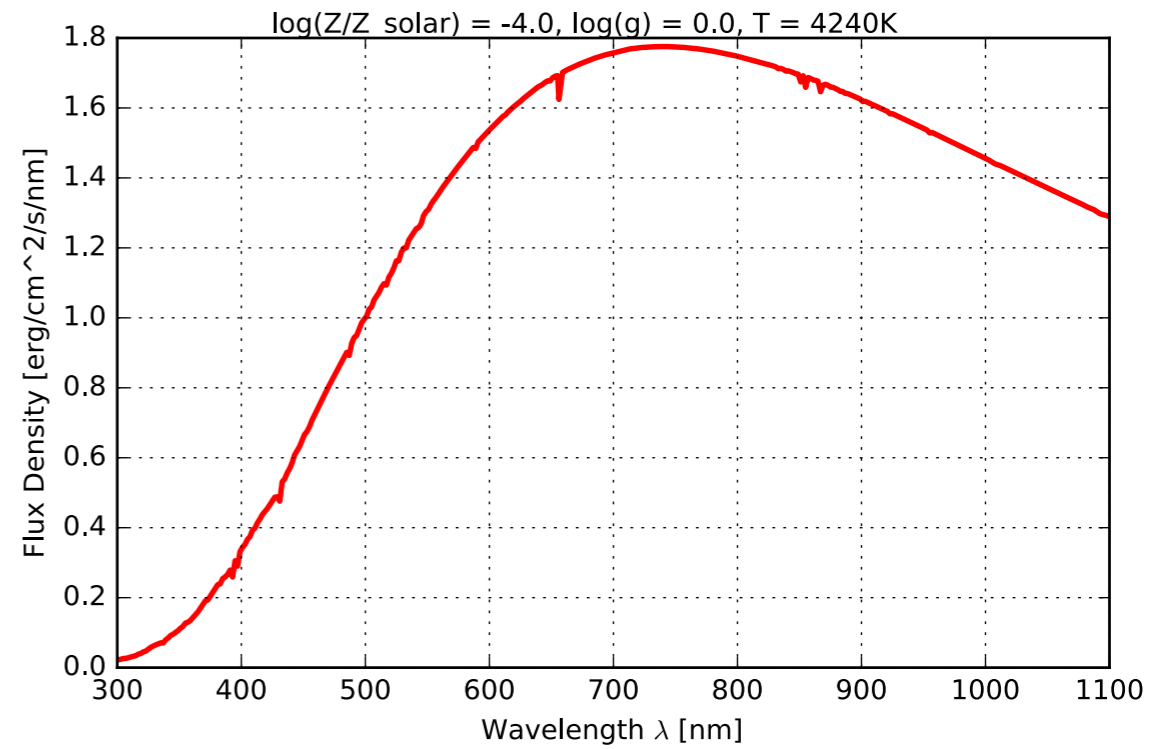
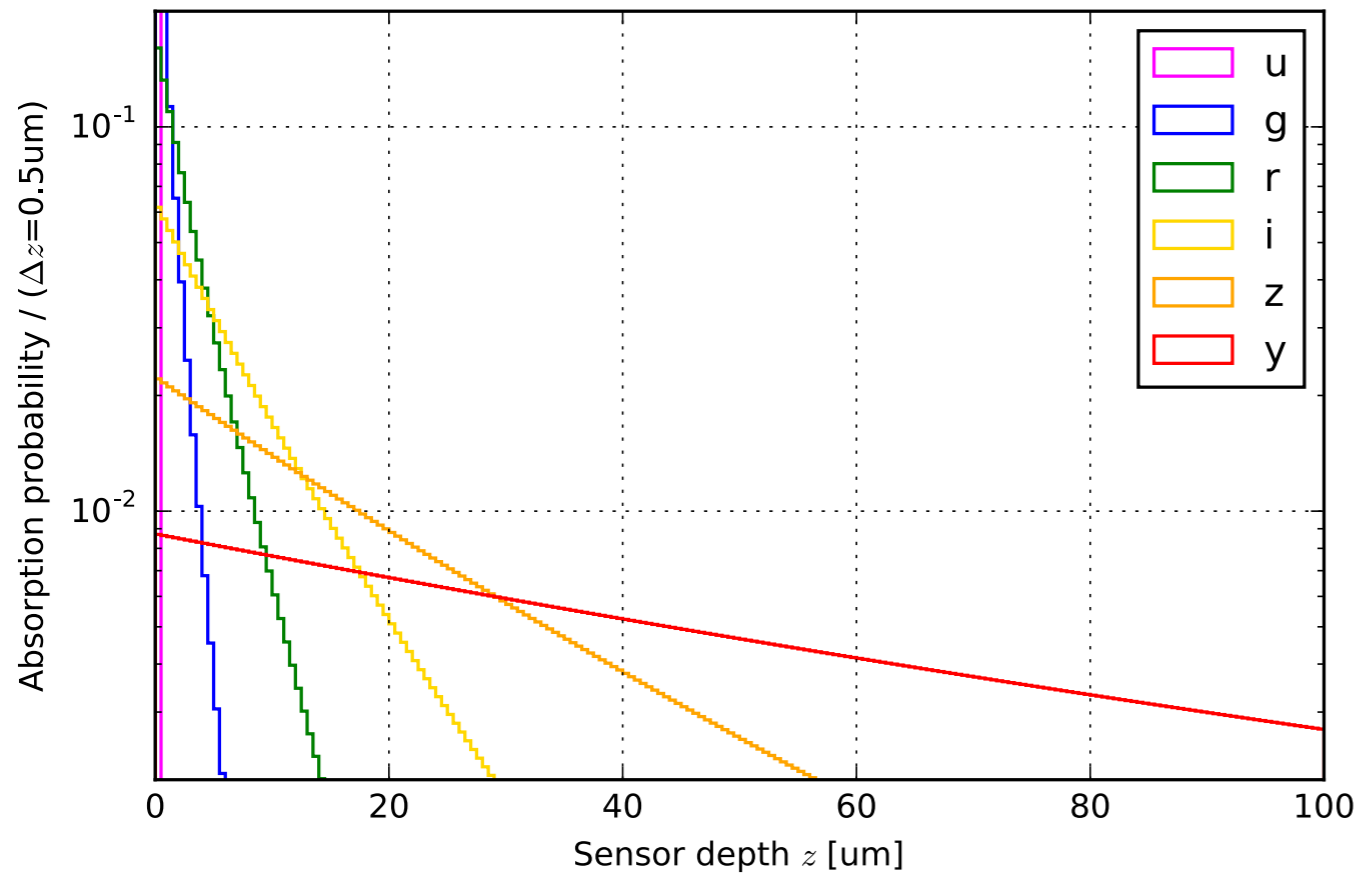
Transmission probability at 90° incidence:

$$R(\lambda) = \frac{(n(\lambda) - n_0)^2 + (k(\lambda) - k_0)^2}{(n(\lambda) + n_0)^2 + (k(\lambda) + k_0)^2}$$



Depth profile of electrons due to a *source* viewed through a *filter*:

$$\frac{1}{N} \frac{dN}{dz} = \int s(\lambda) f(\lambda) \frac{e^{-z/a(\lambda)}}{A(\lambda, \Delta z)} d\lambda$$



How does “brighter-fatter” effect vary with the filter?

Solve Poisson’s equation in quasi-static limit:

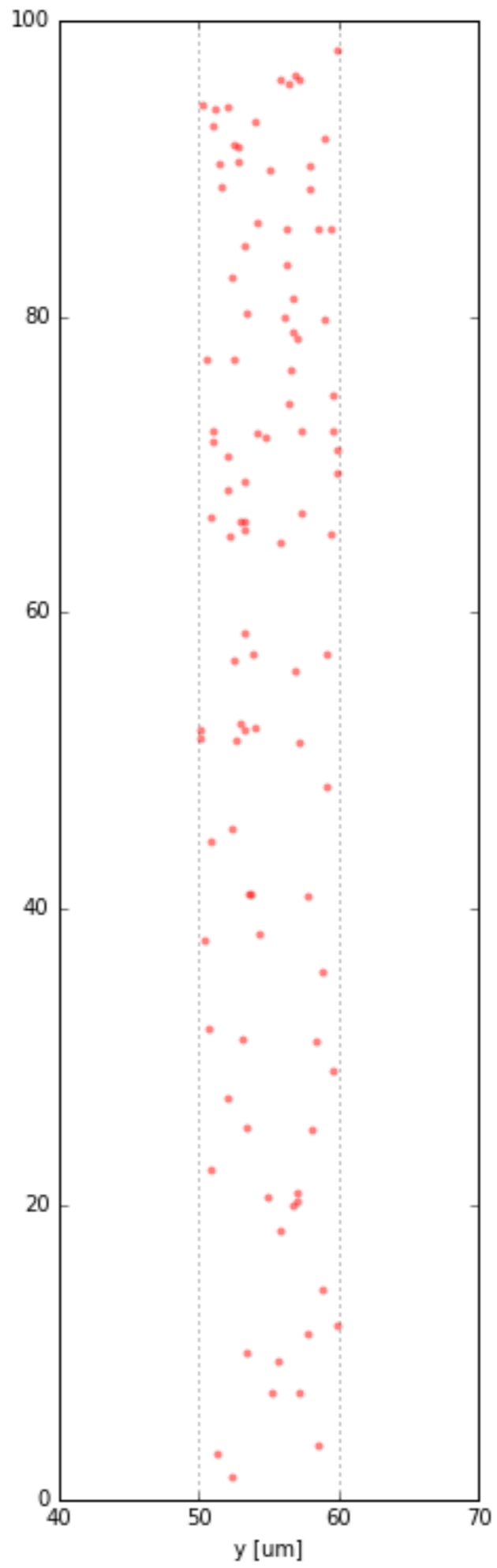
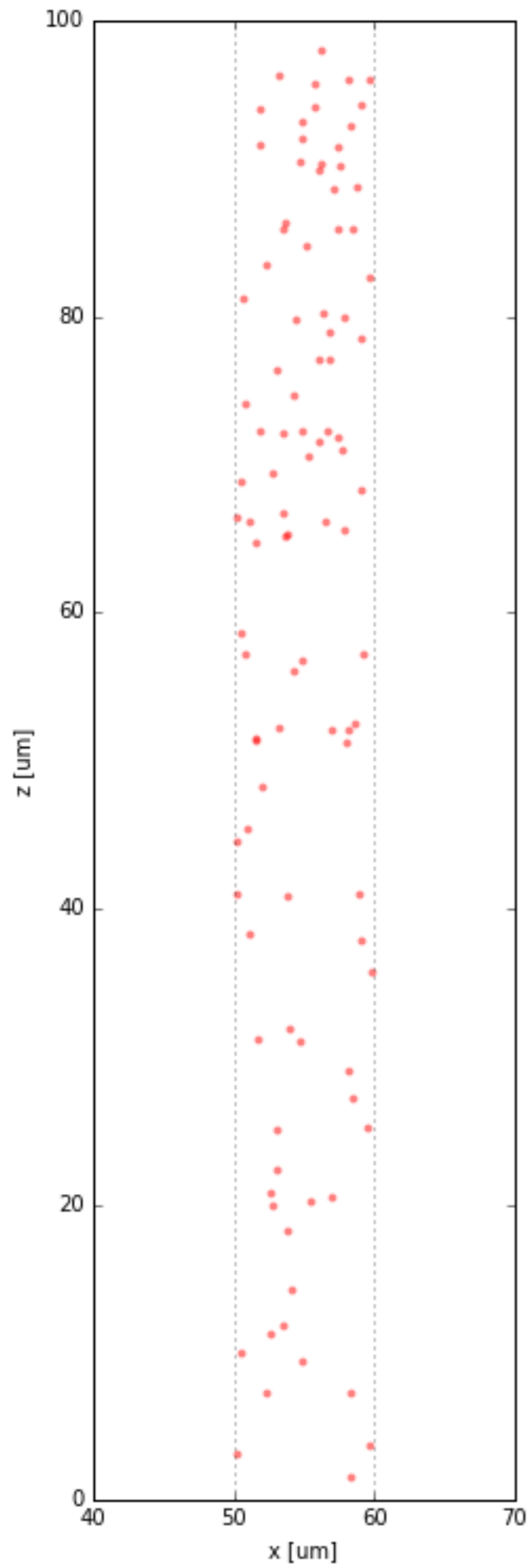
$$\nabla \cdot (\epsilon(x) \nabla \phi(x)) = \rho(x)$$

*permittivity
(geometry &
materials)*

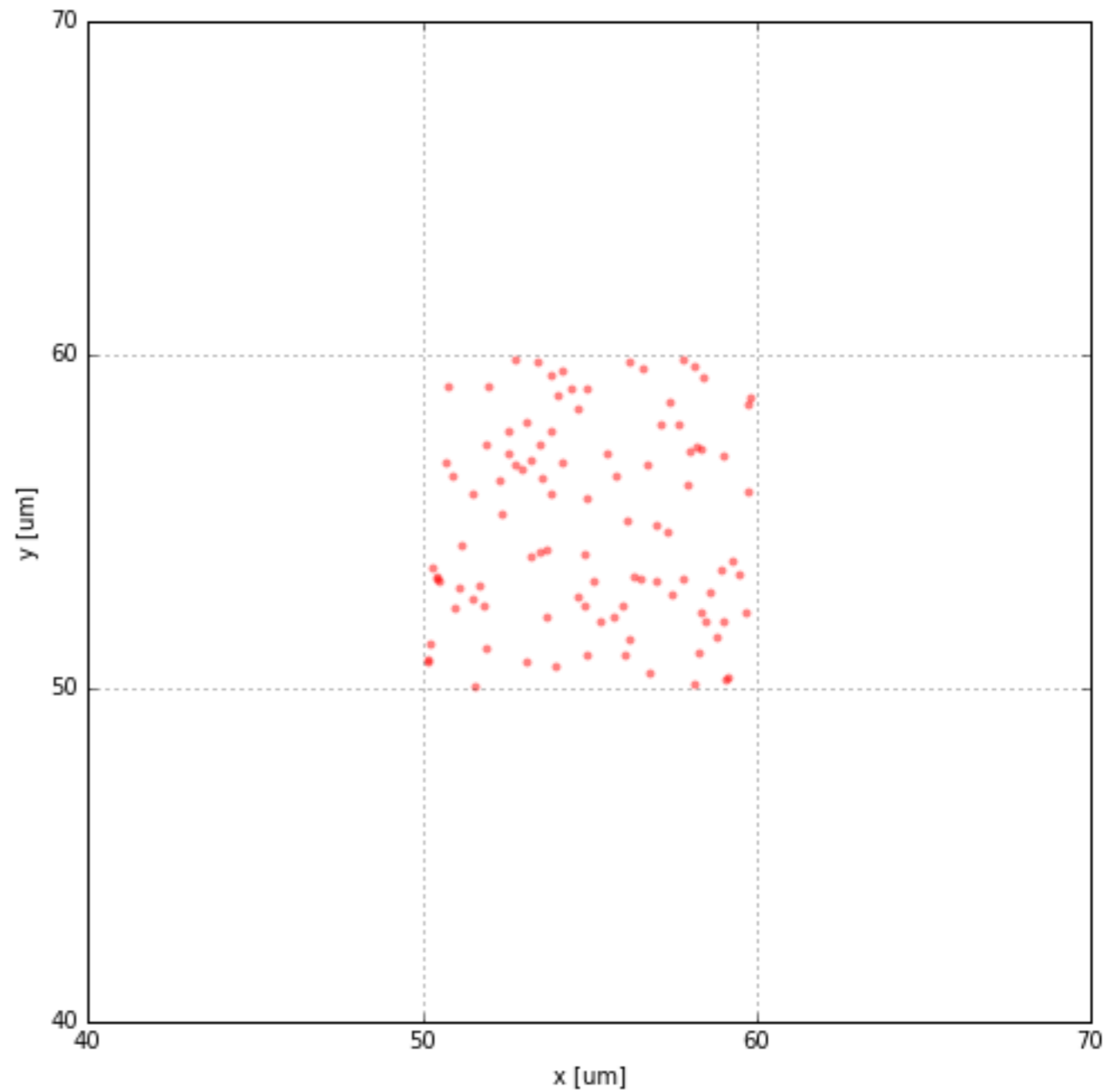
*potential
(bias & gate
voltages)*

*charge density
(static ions &
mobile $e + h$)*

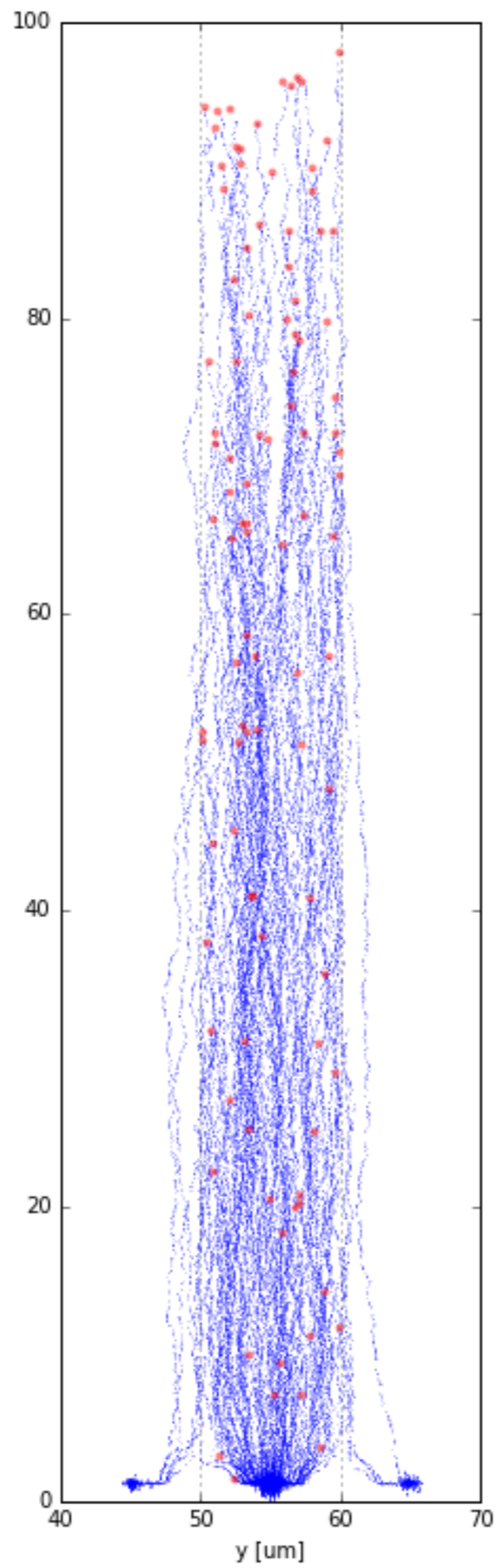
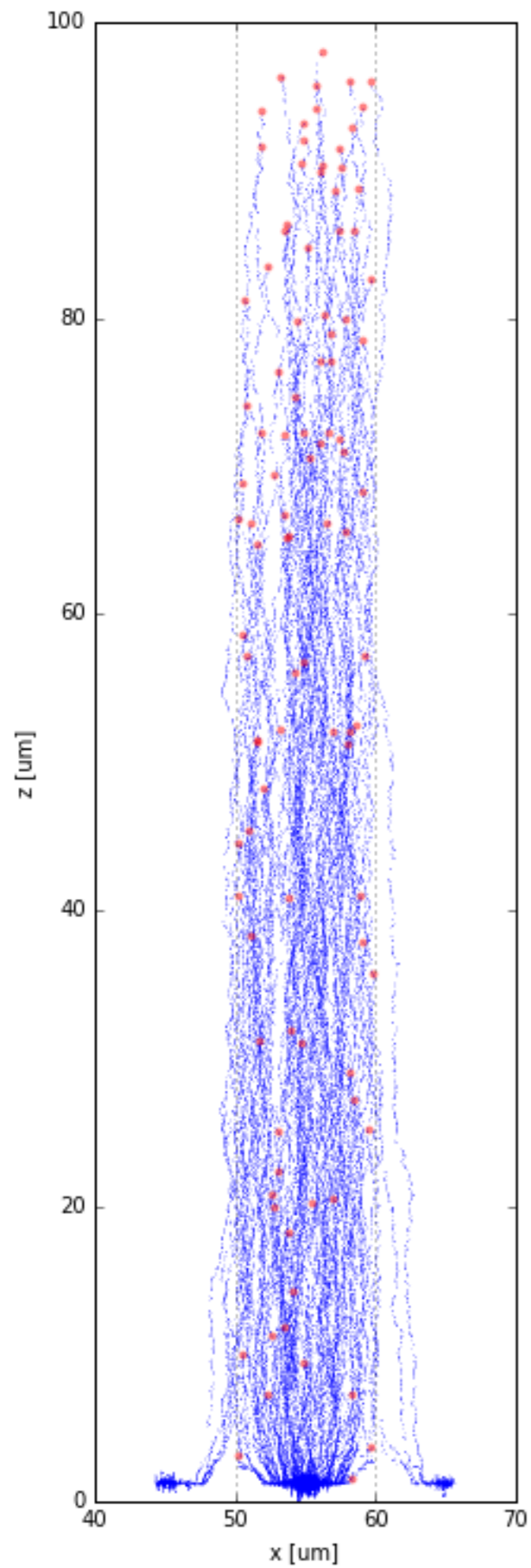
(Implemented in the `wavelength_effects` branch of `Poisson_CCD22` repo.)



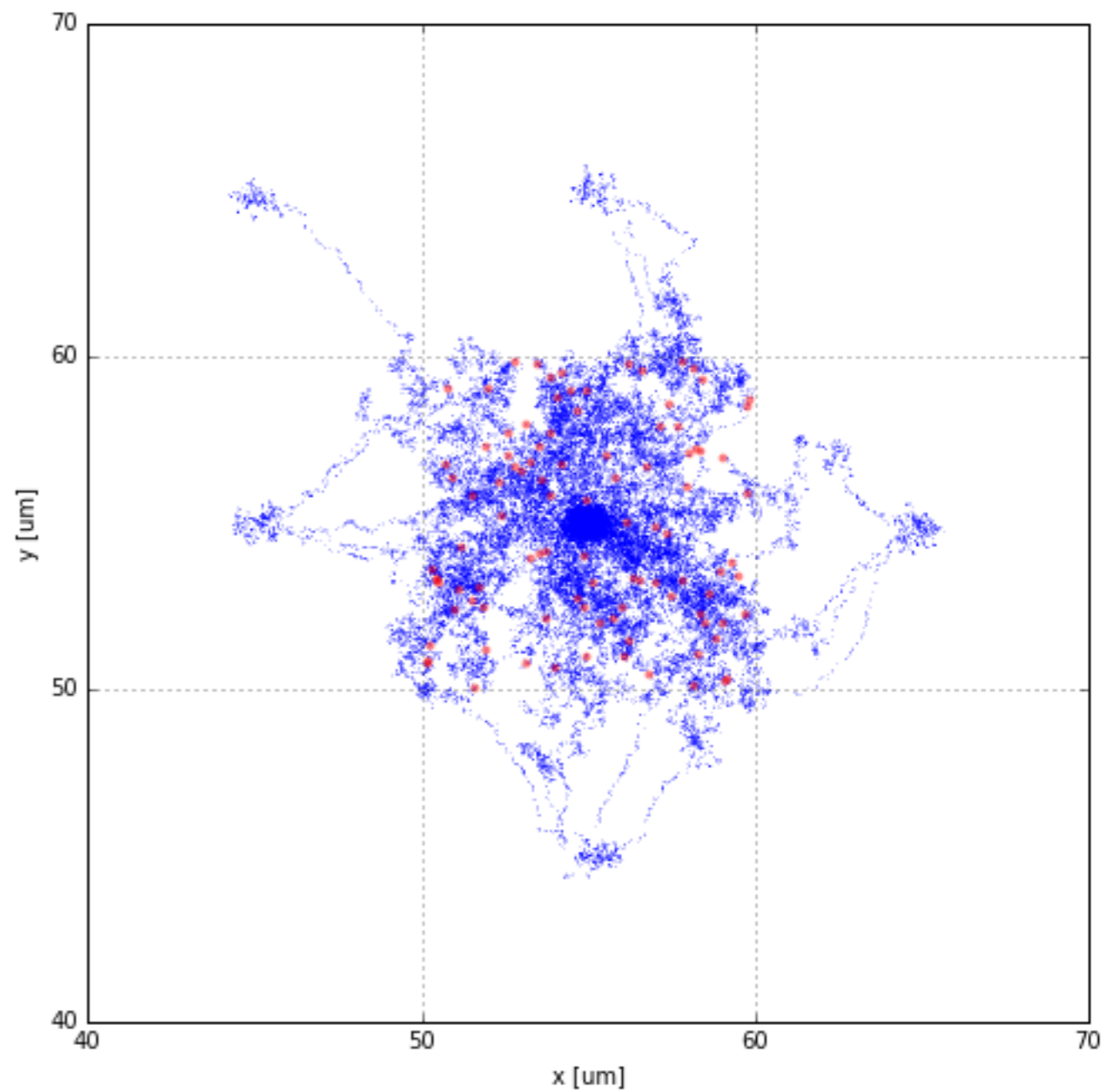
Y BAND
NO ACCUMULATED SIGNAL



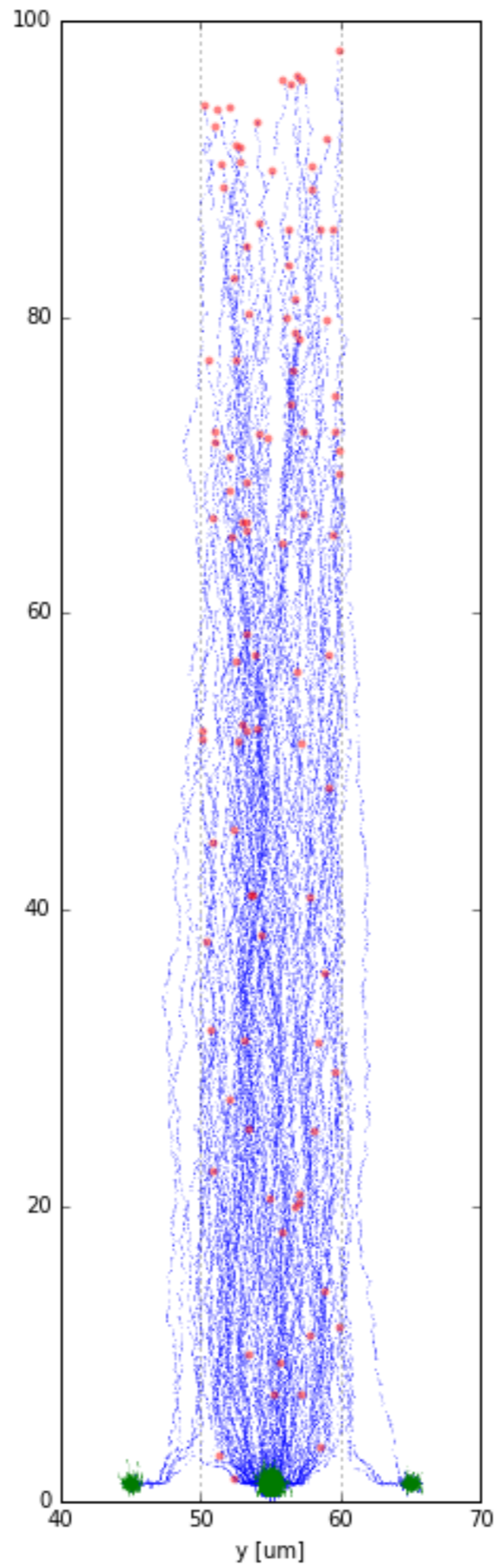
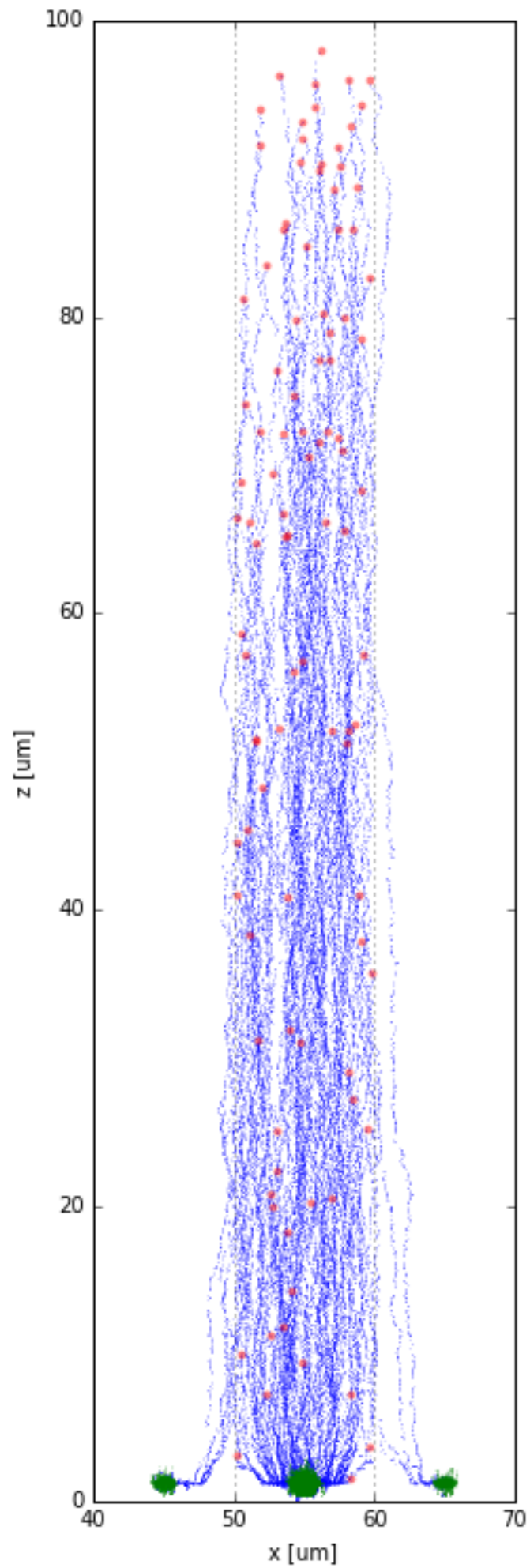
Initial electron positions



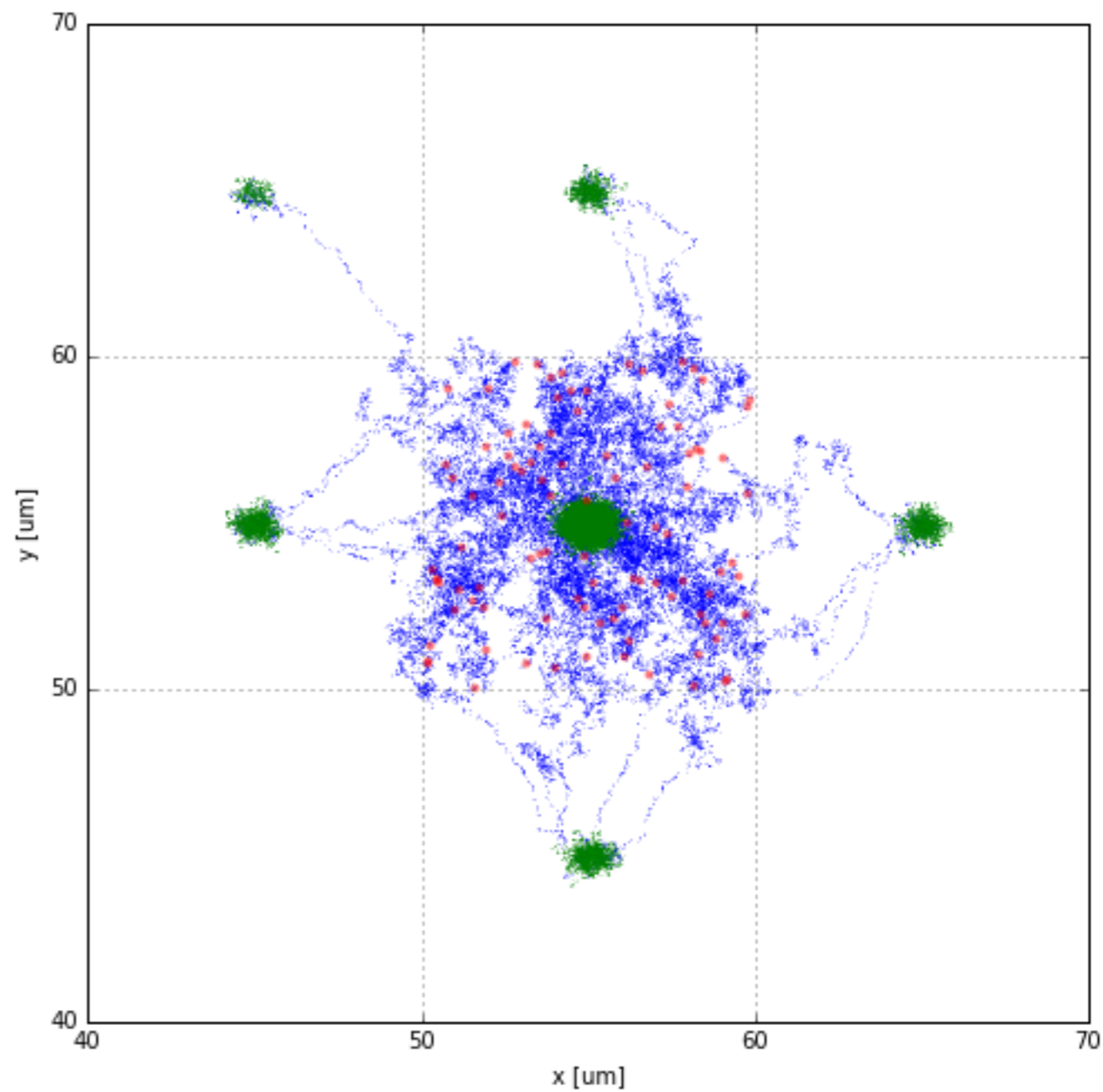
Y BAND
NO ACCUMULATED SIGNAL



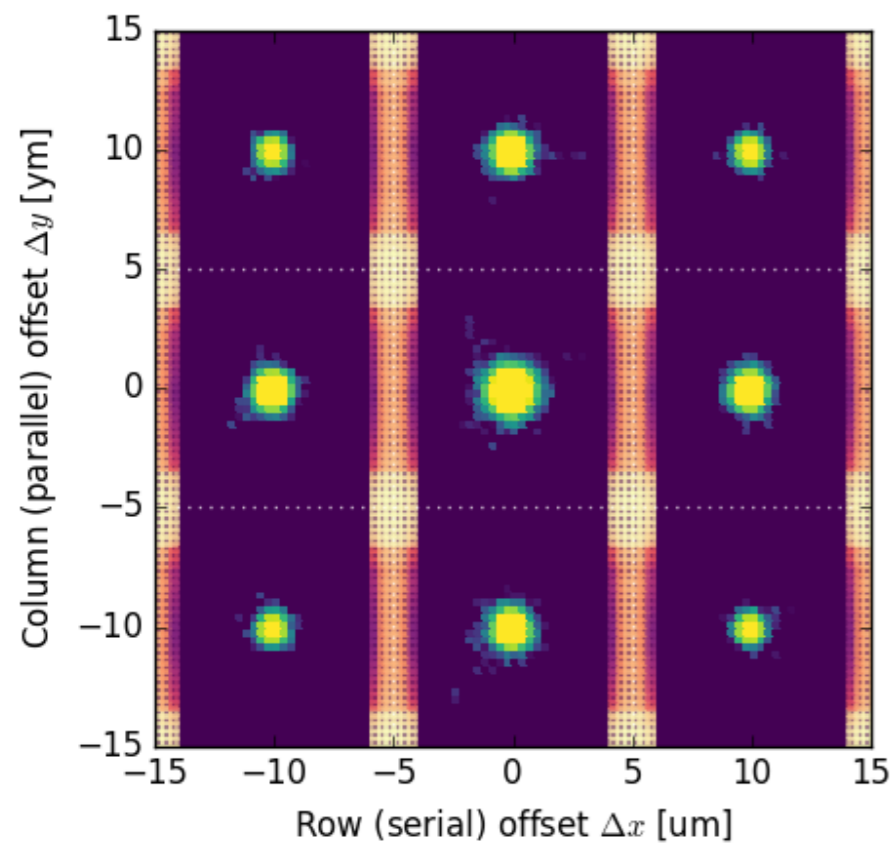
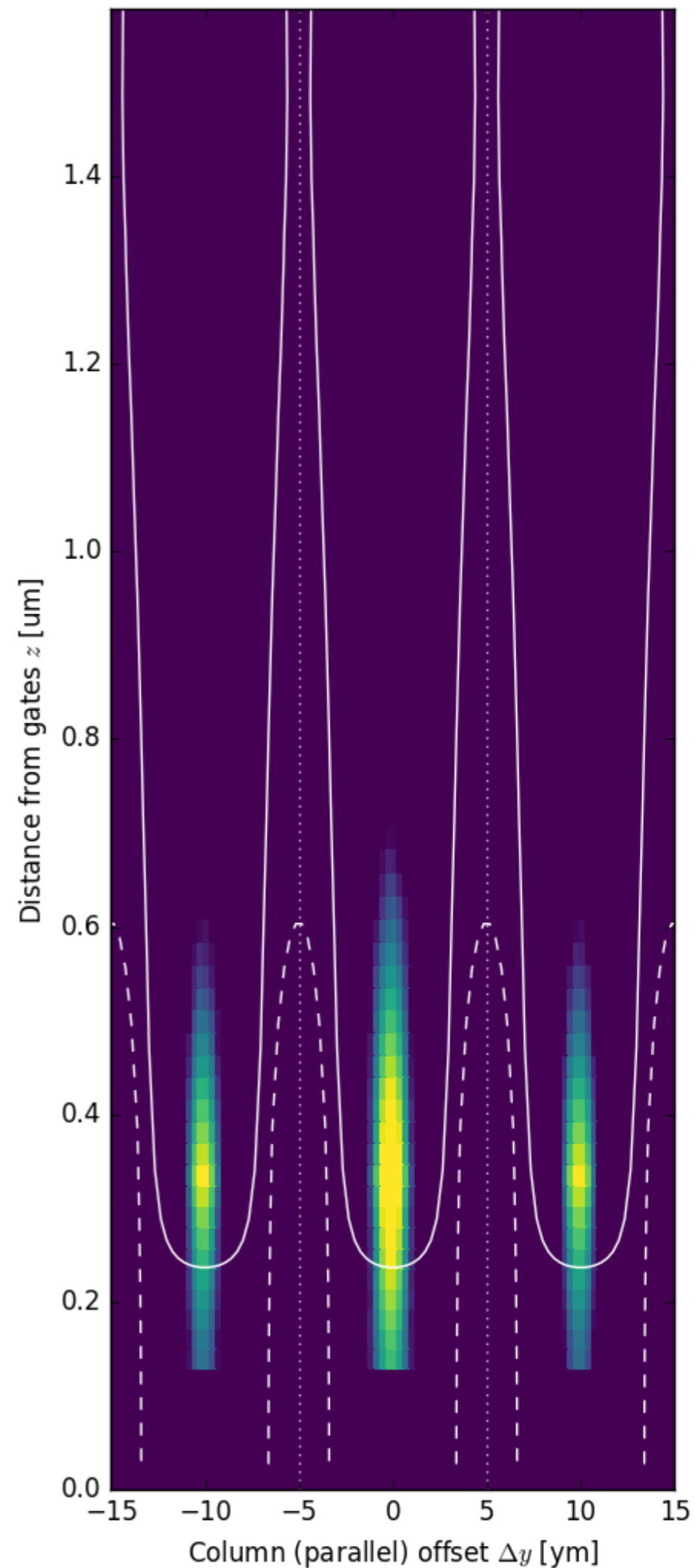
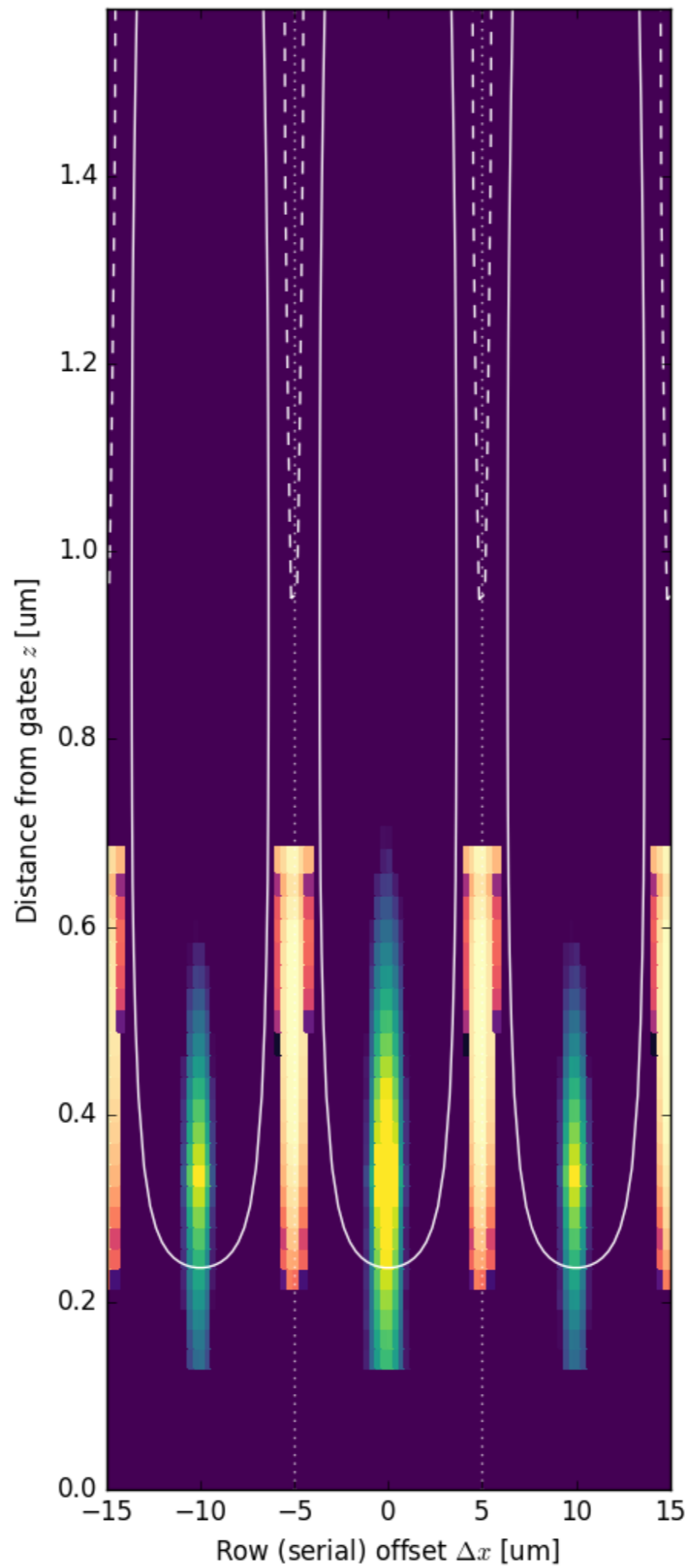
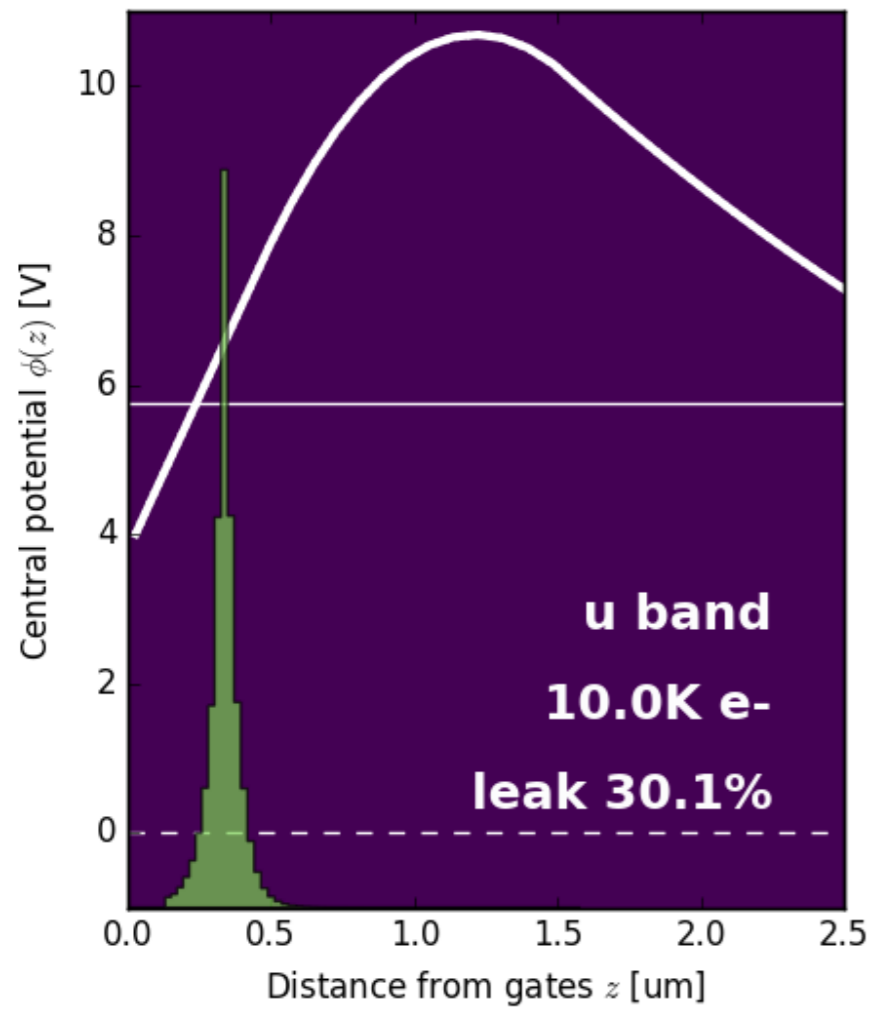
Electron tracks
(diffusion + E-field)

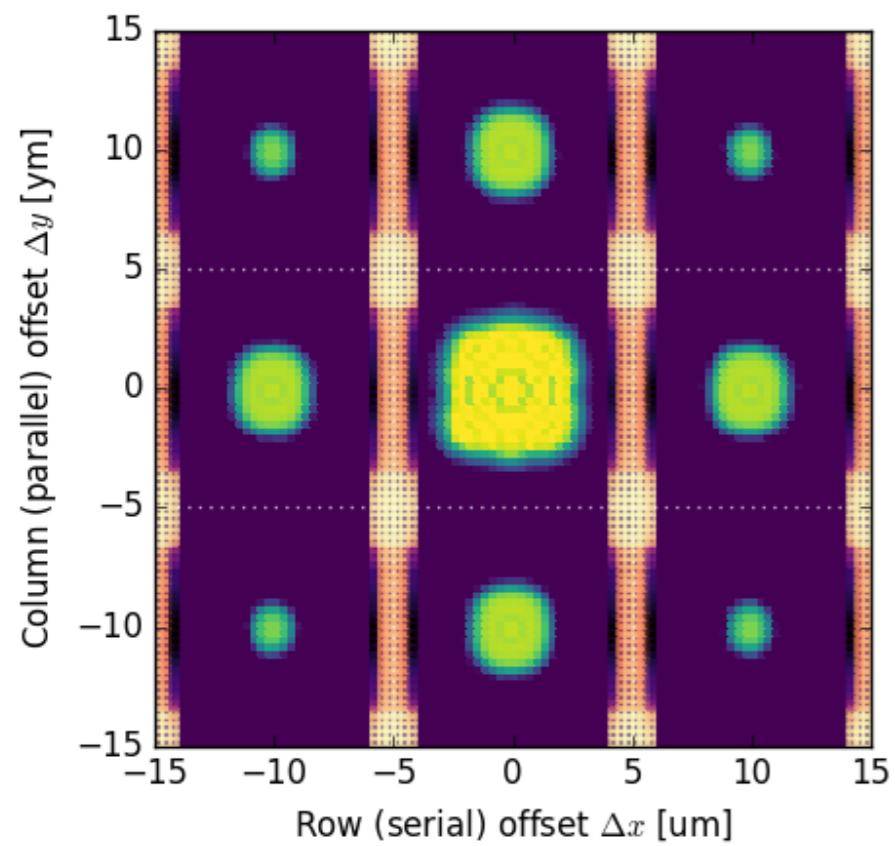
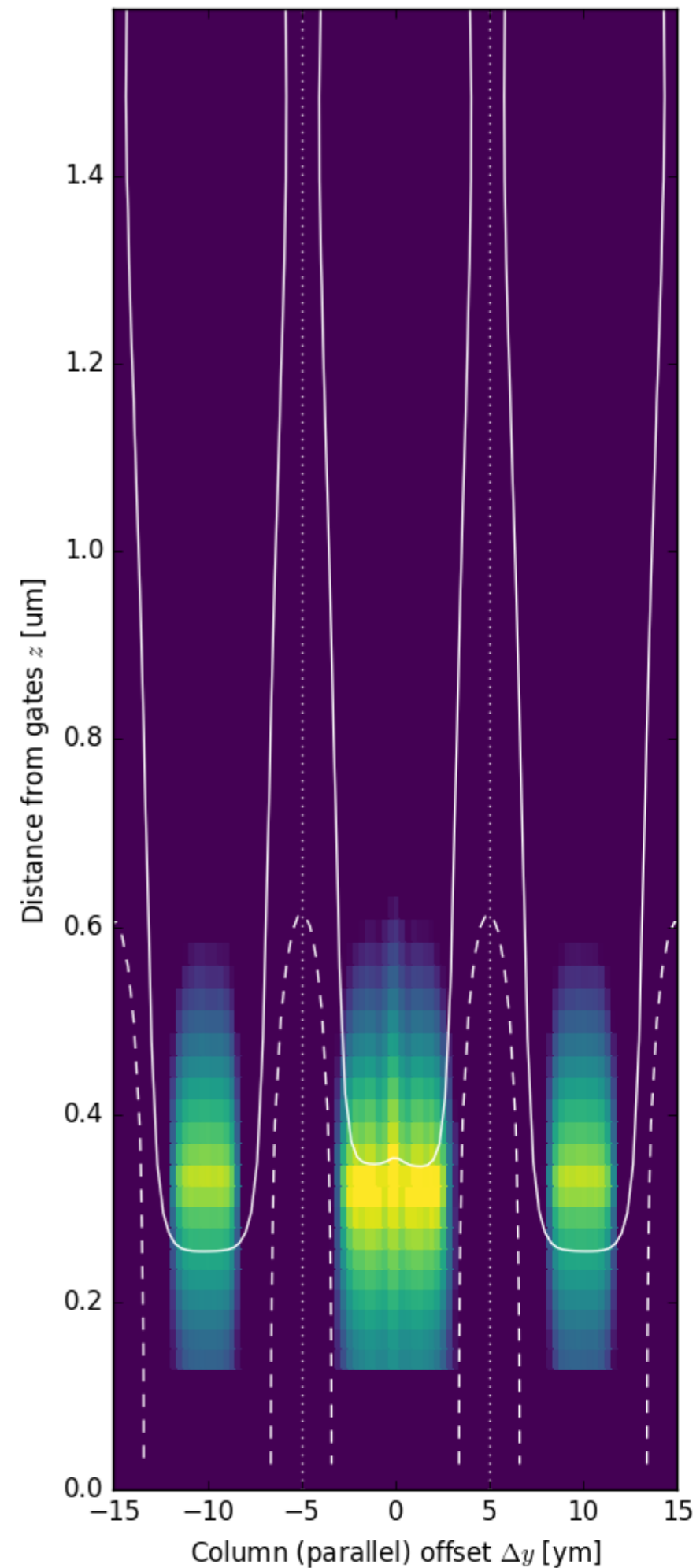
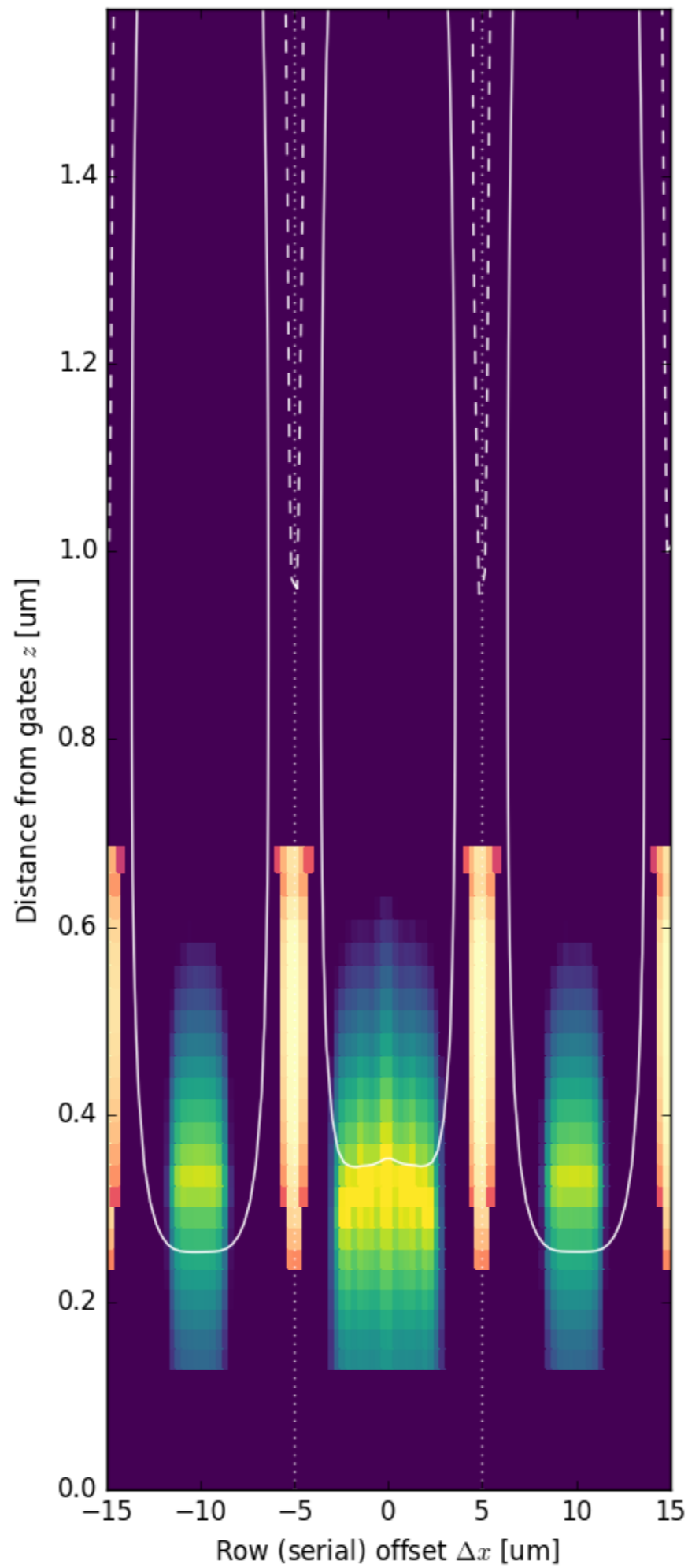
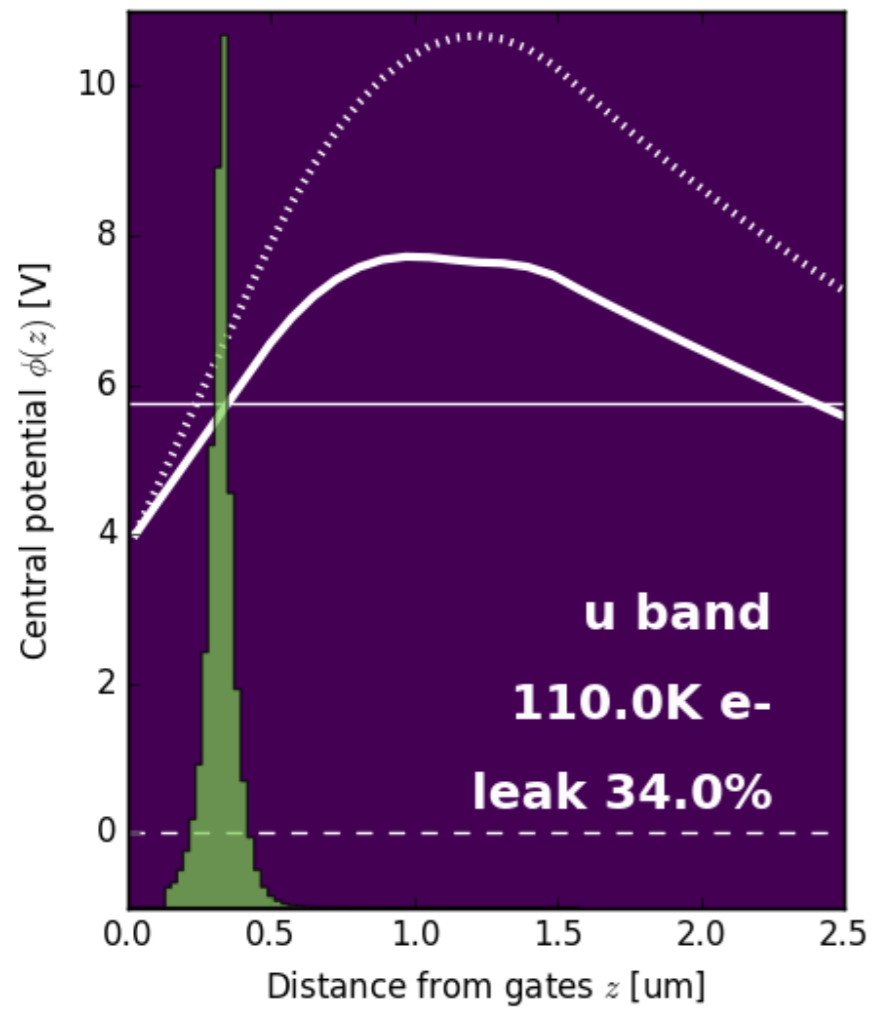


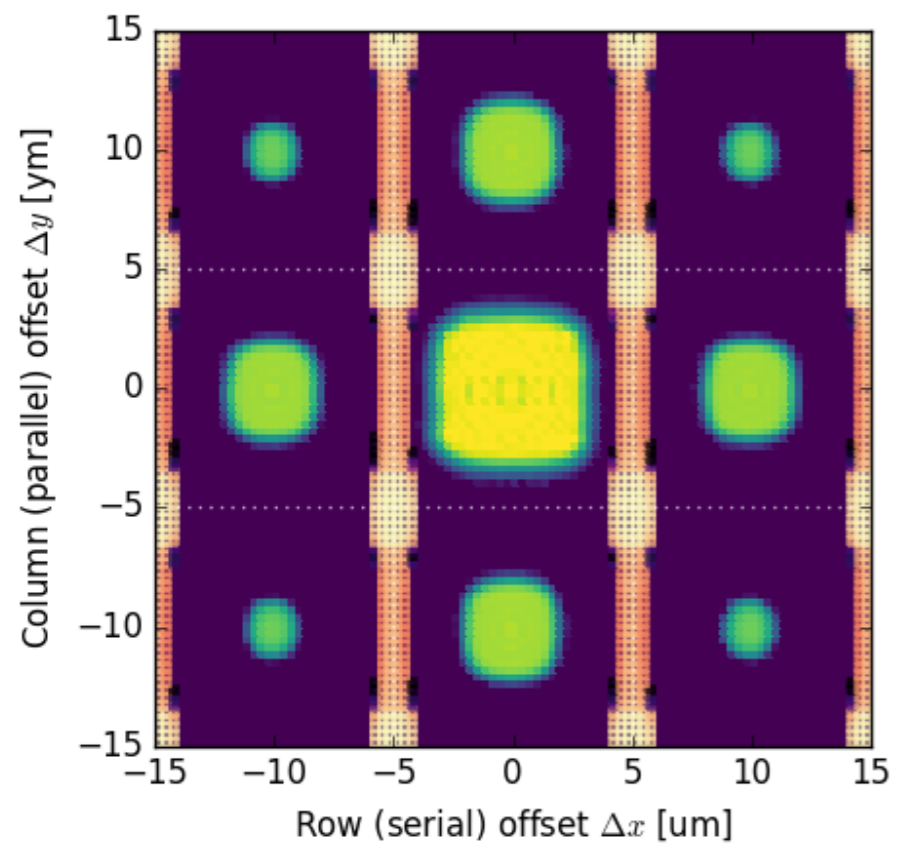
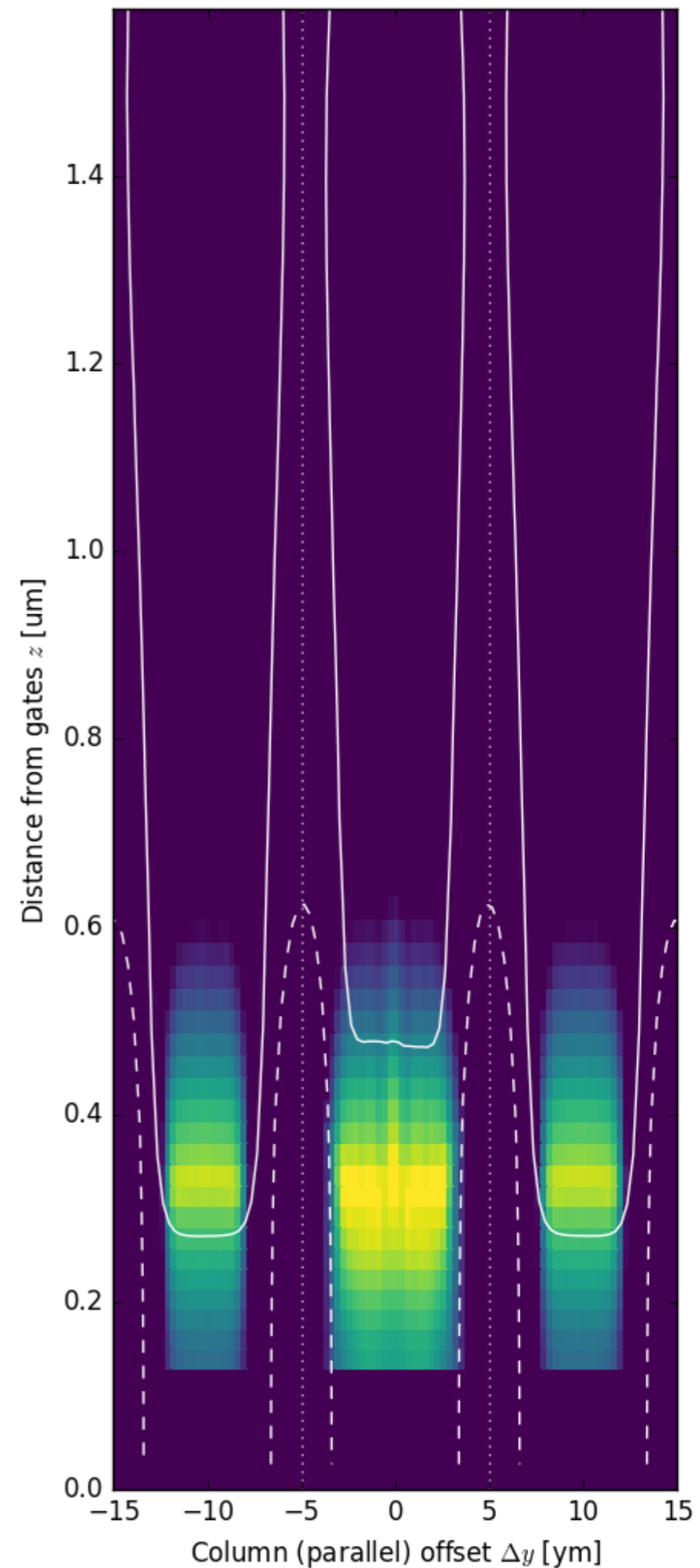
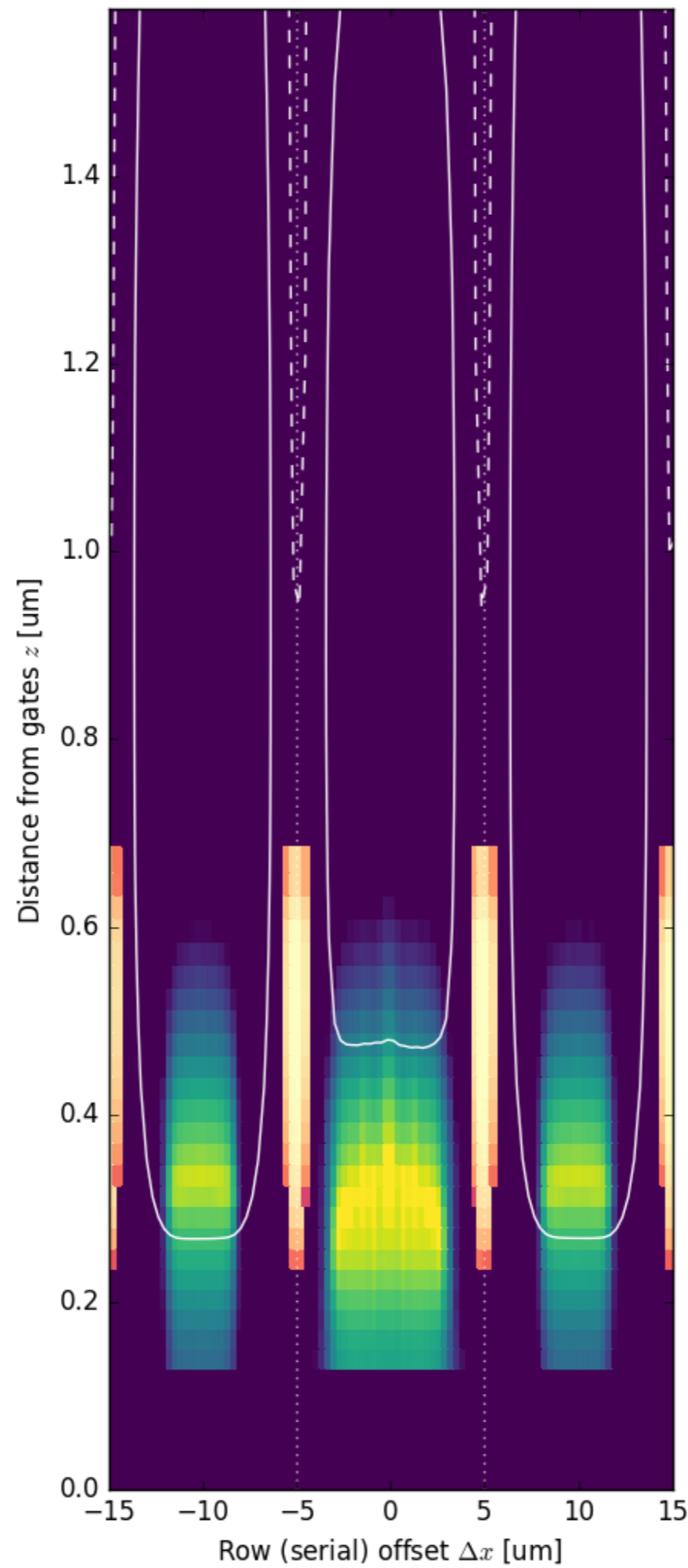
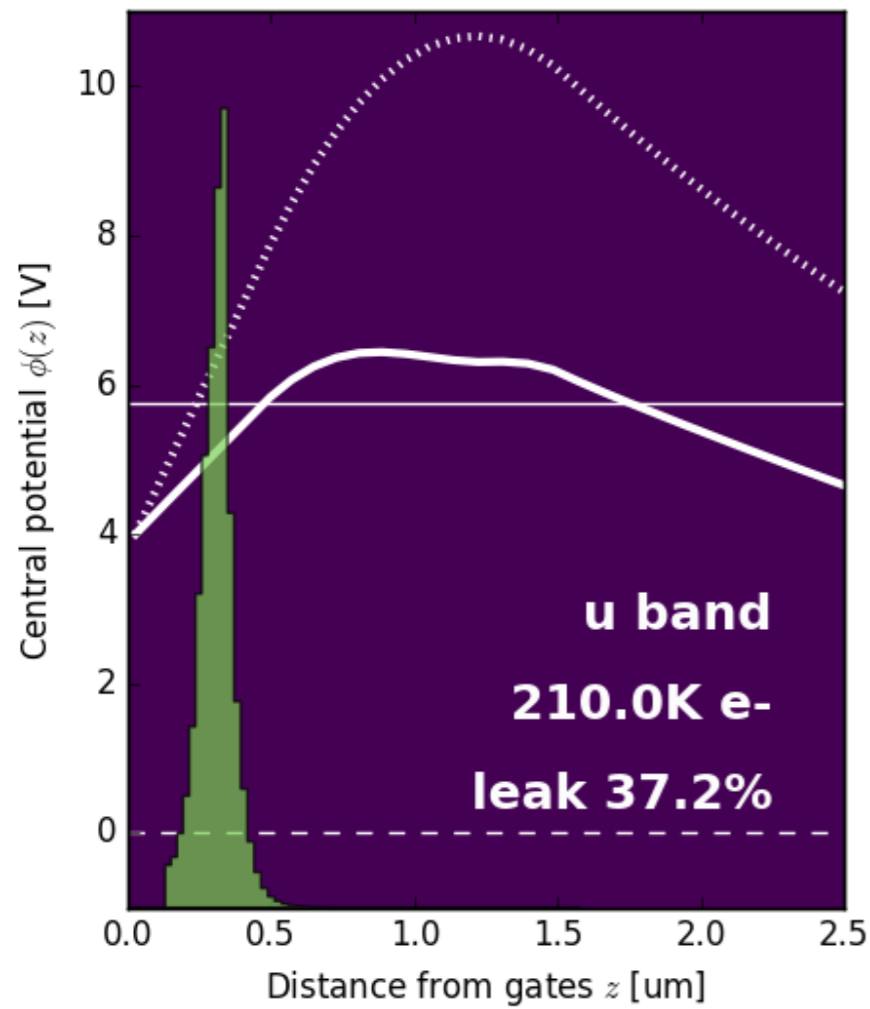
Y BAND
NO ACCUMULATED SIGNAL

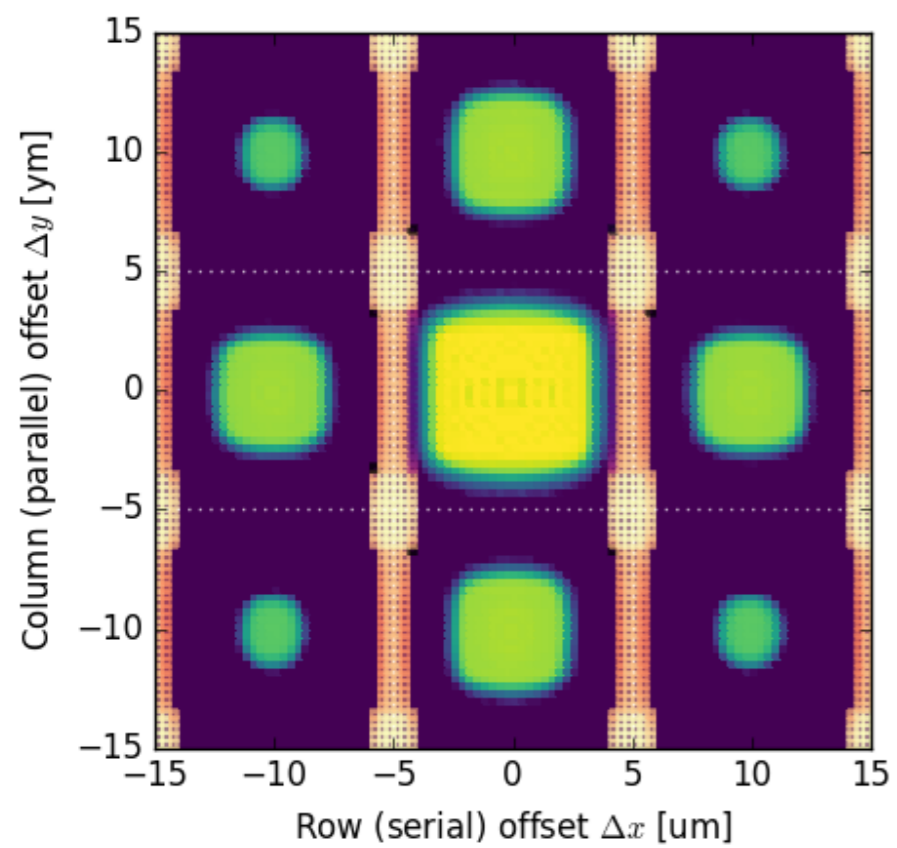
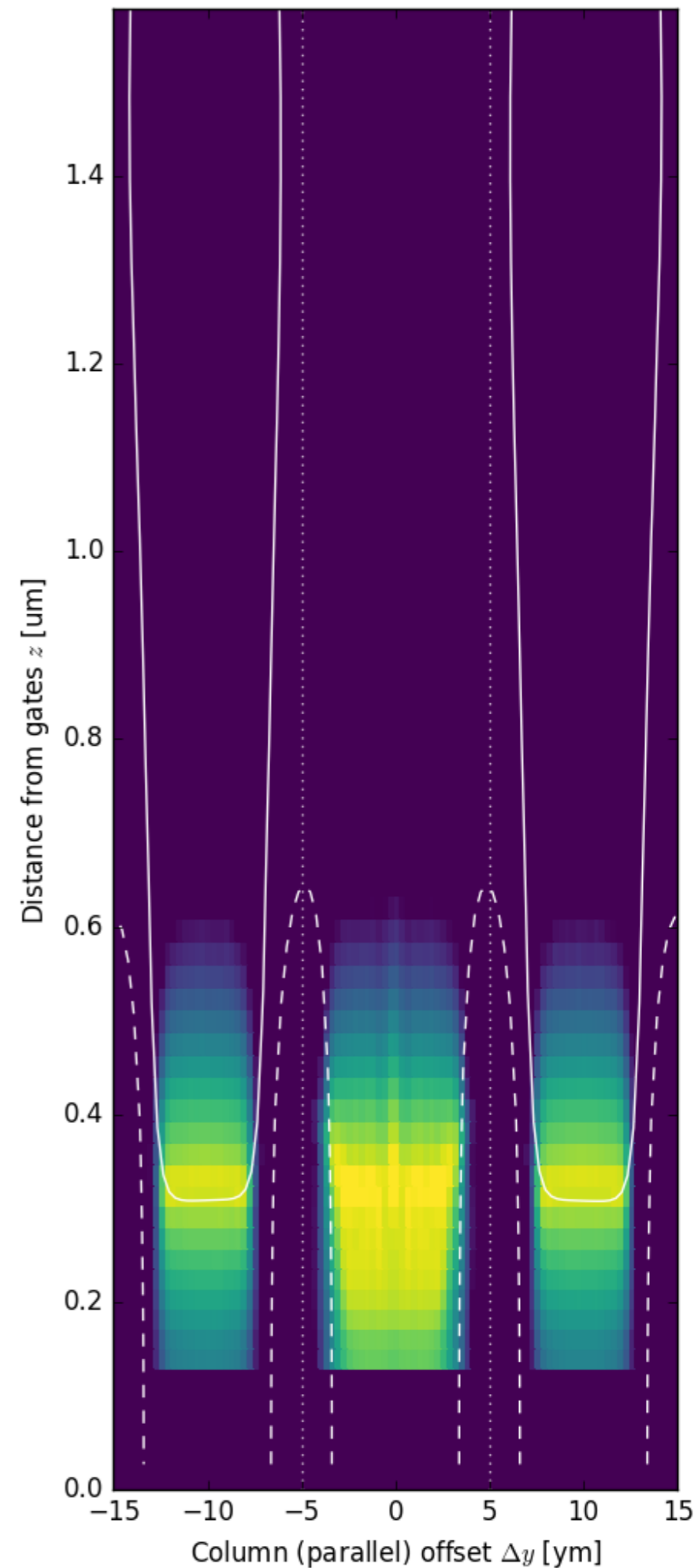
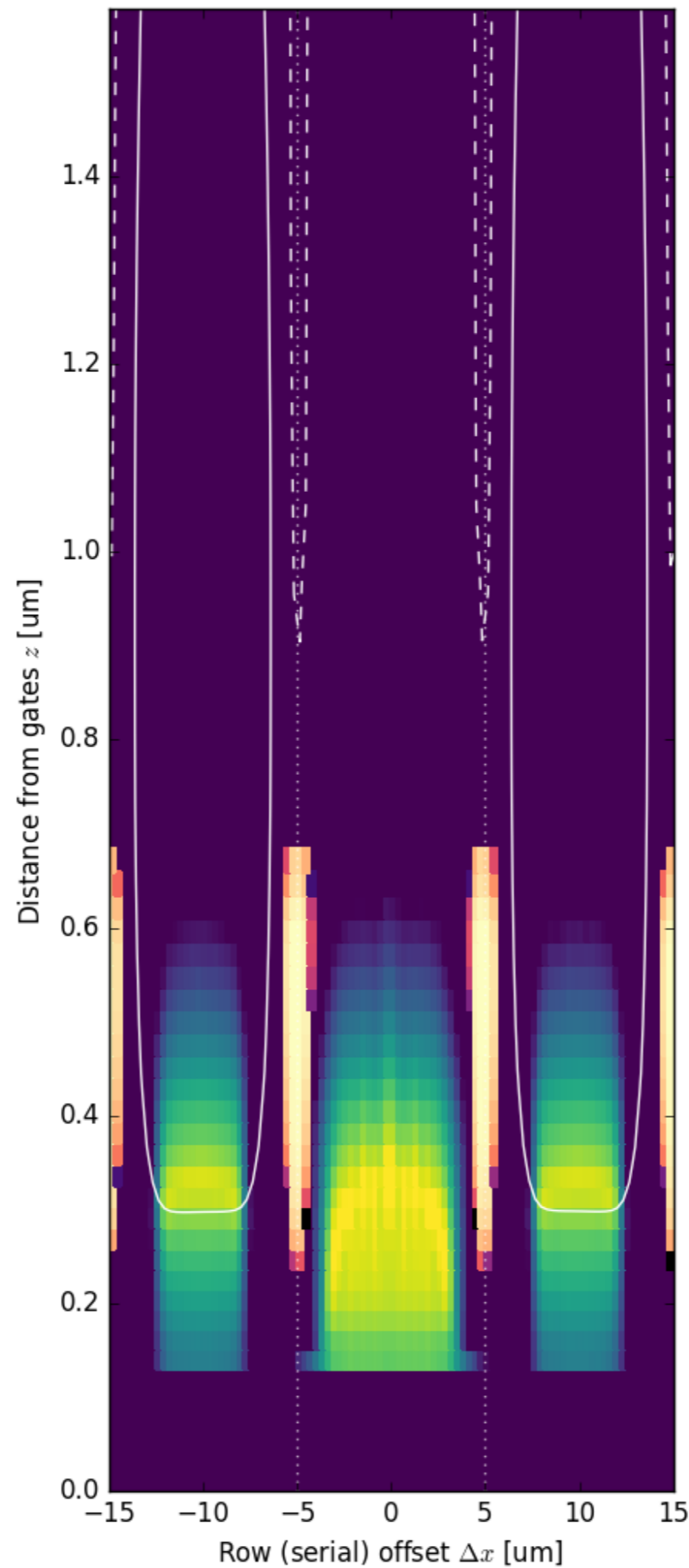
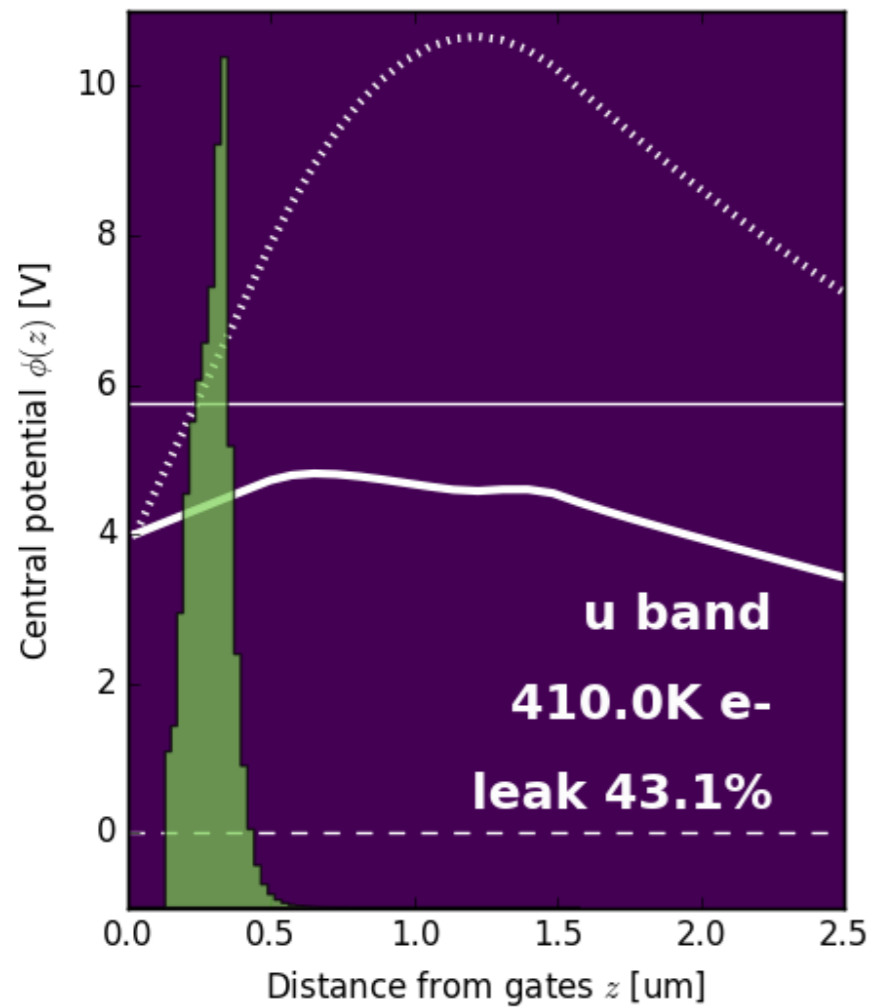


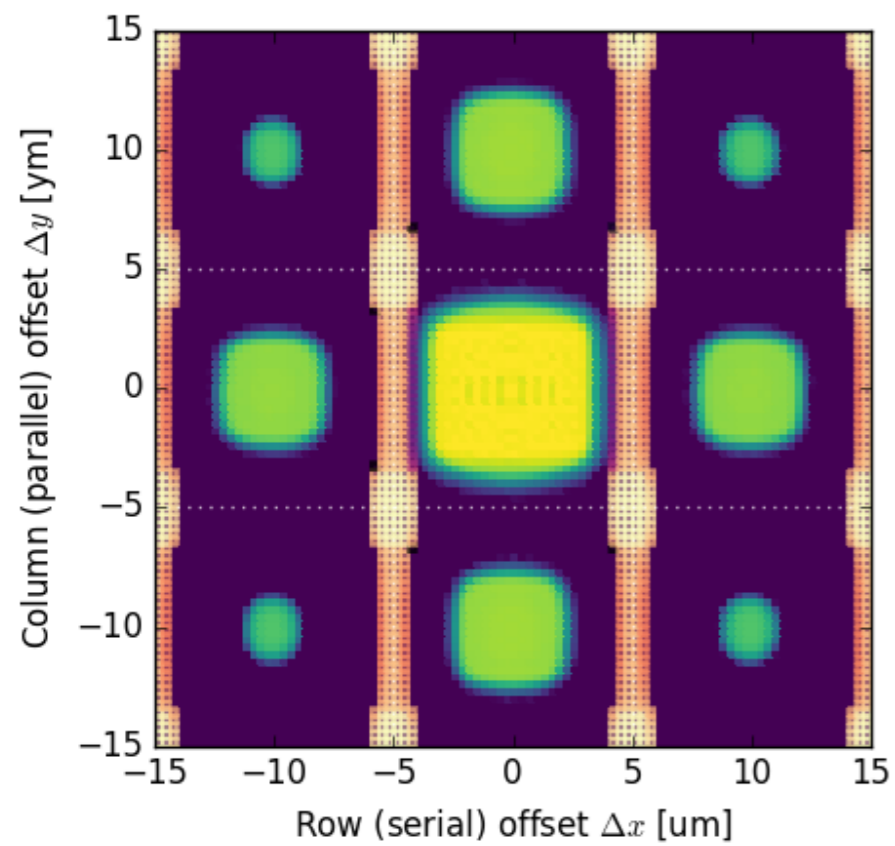
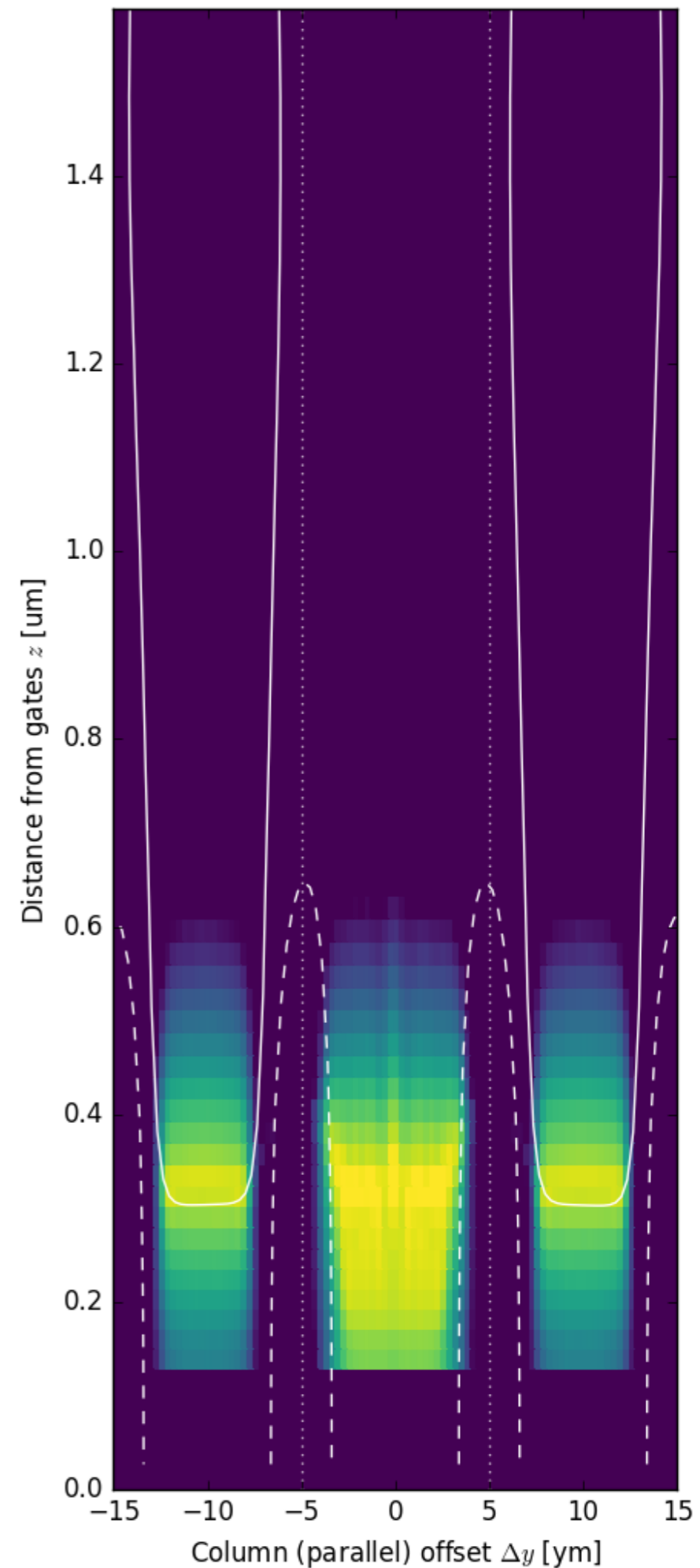
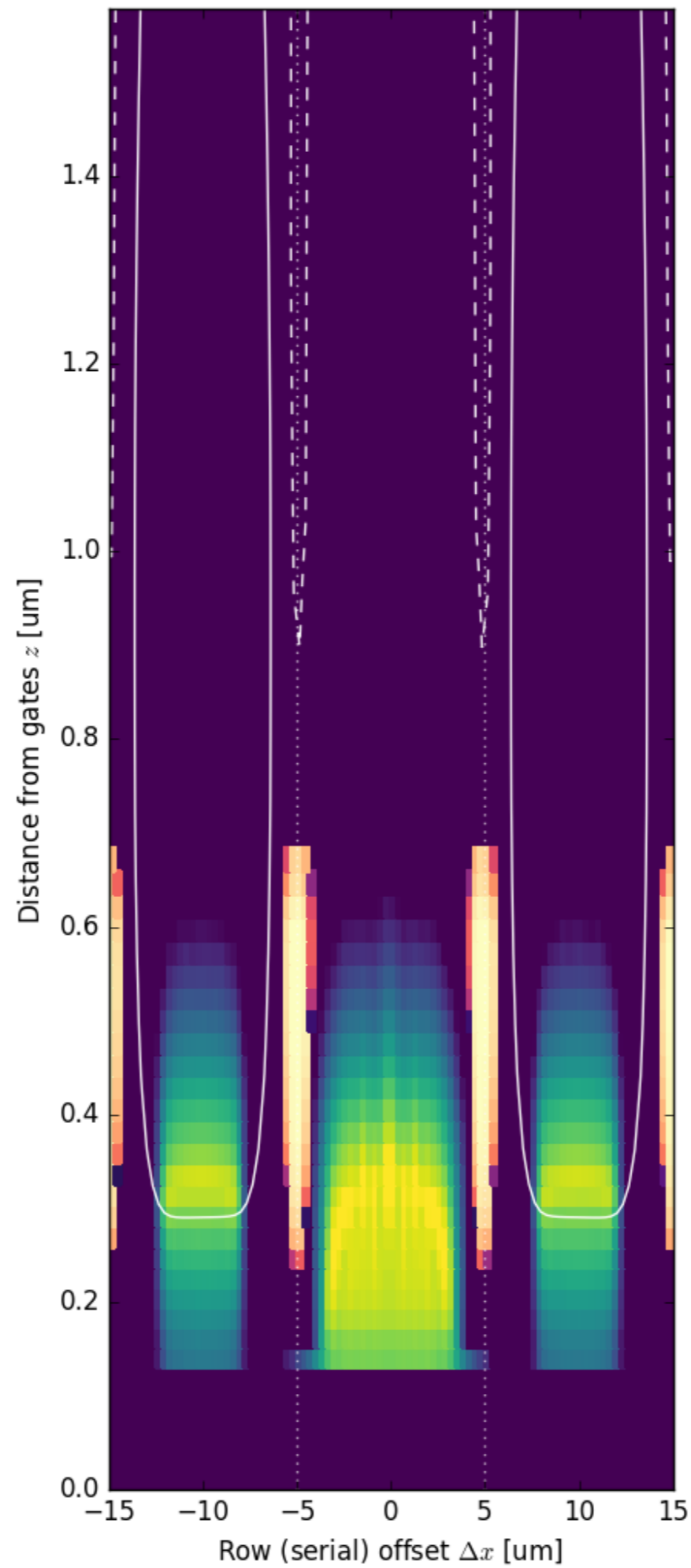
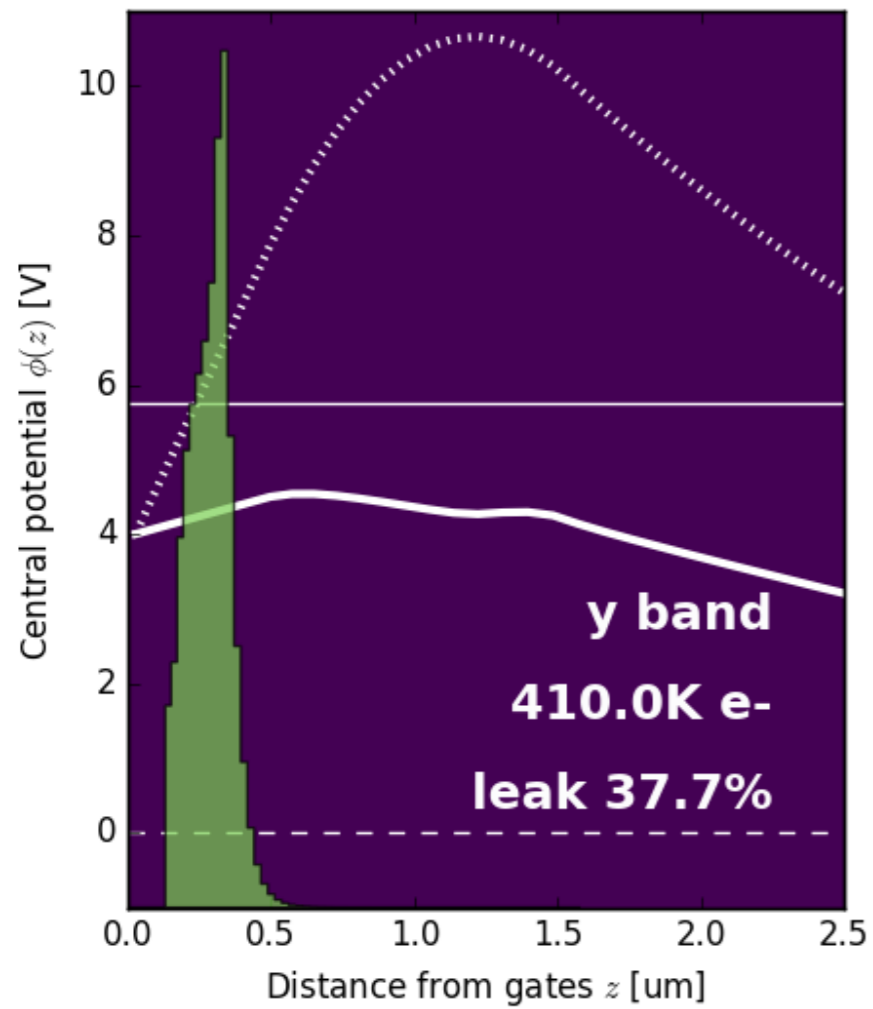
**“Equilibrium” charge
distribution**

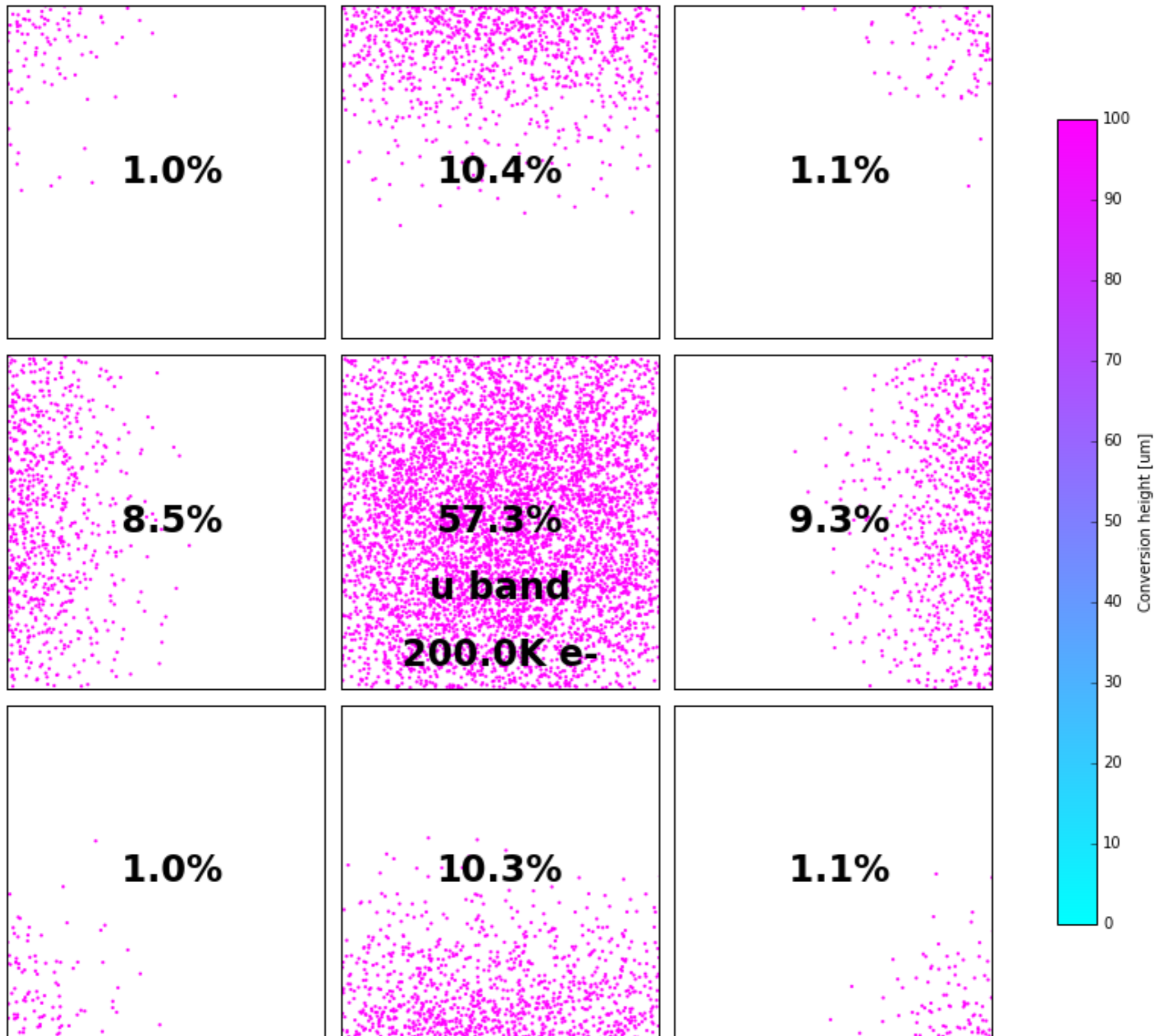


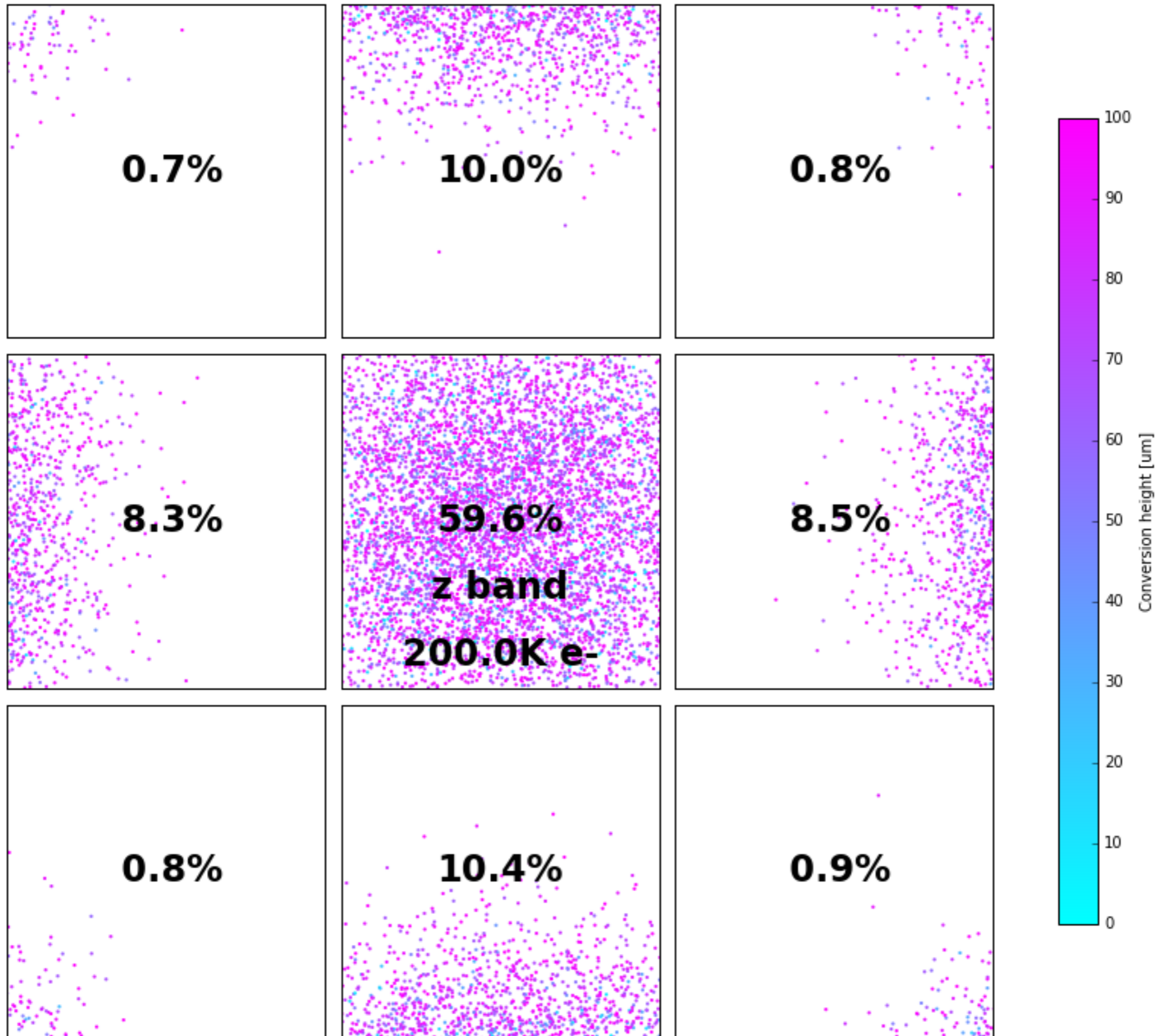


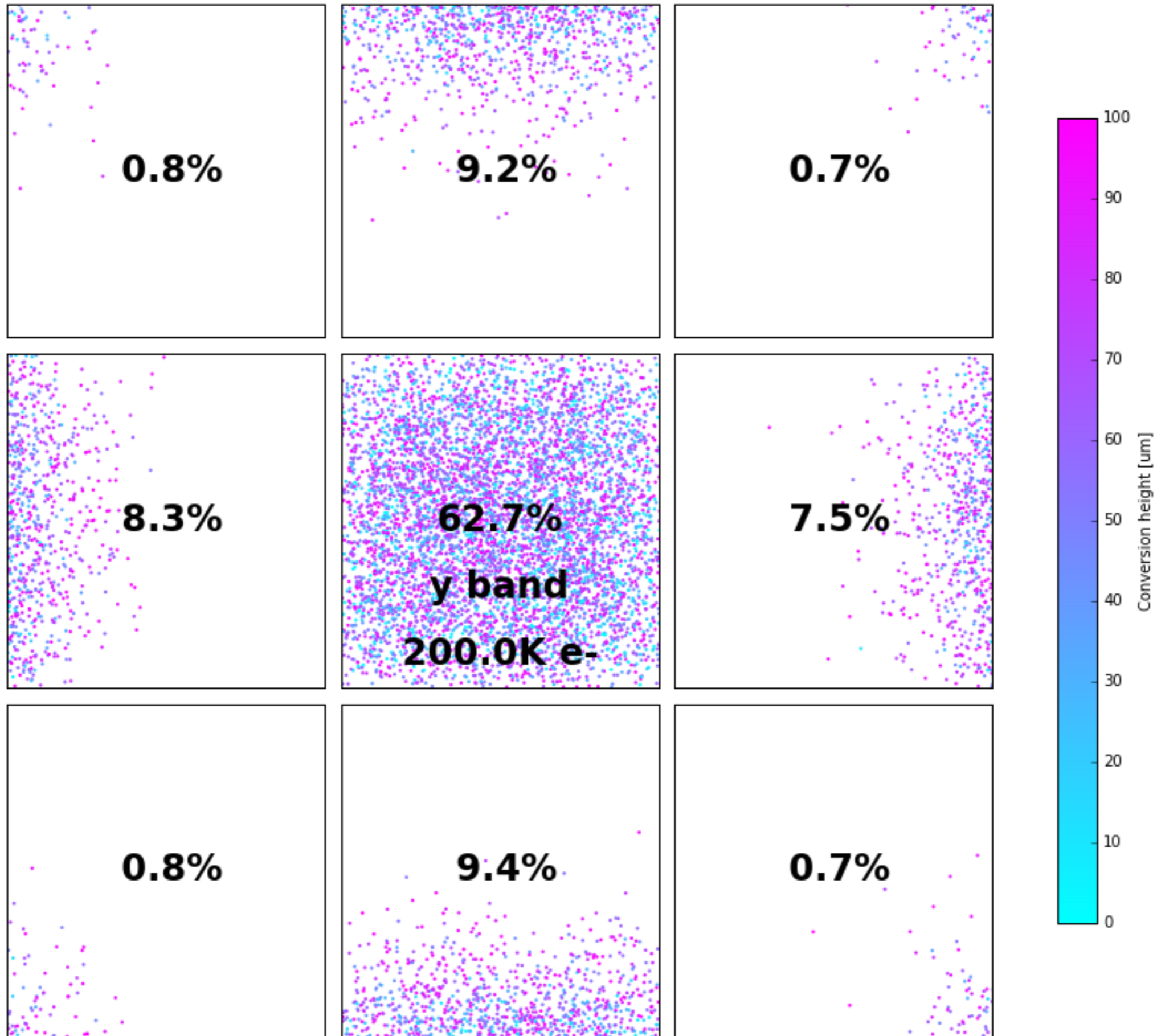


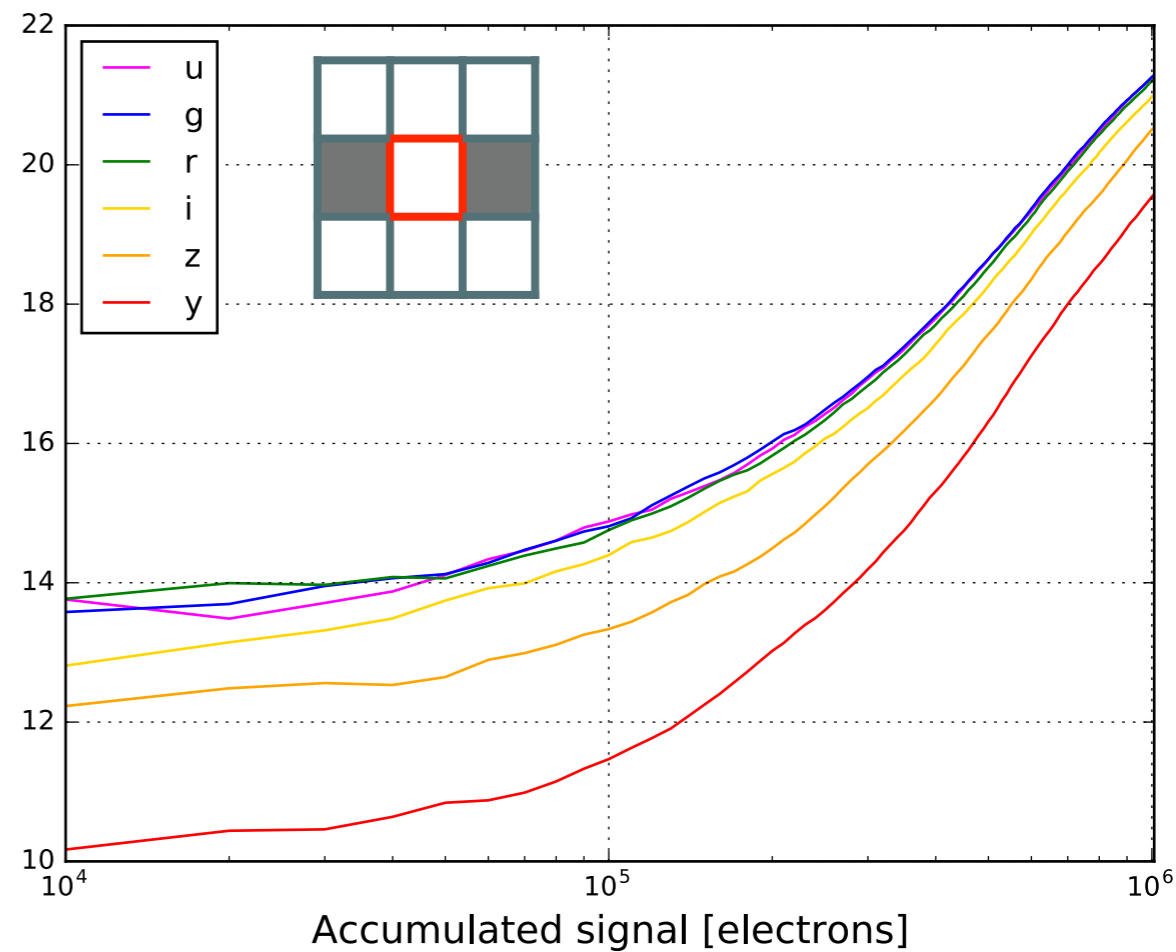
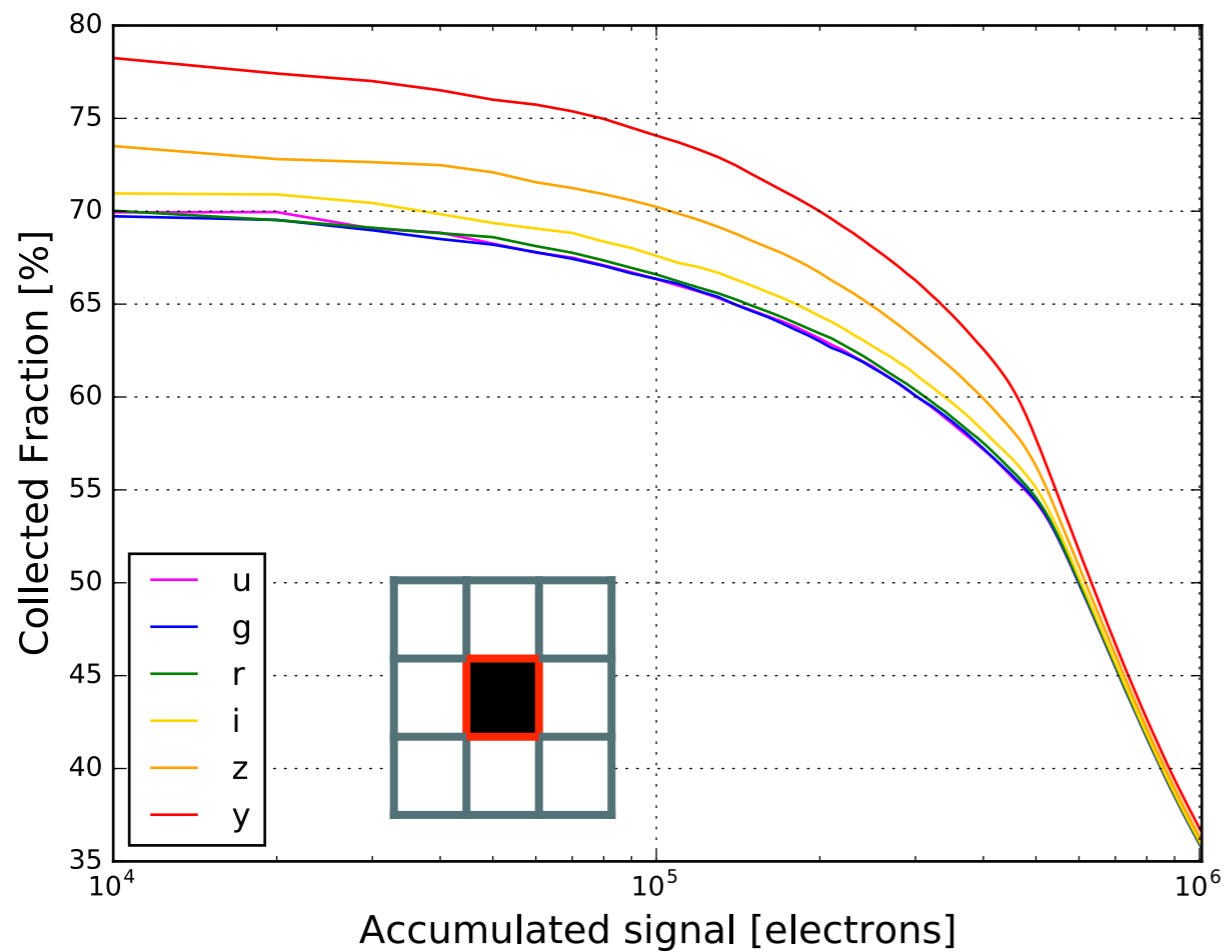
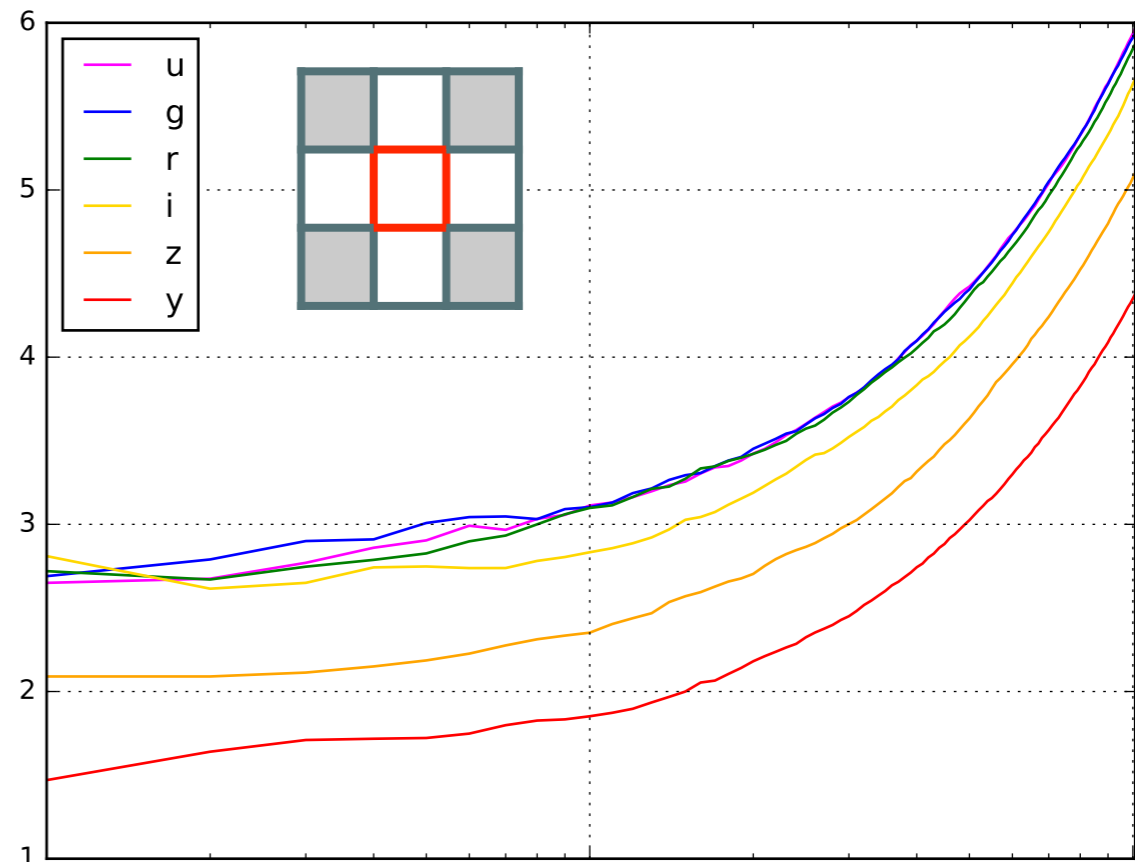
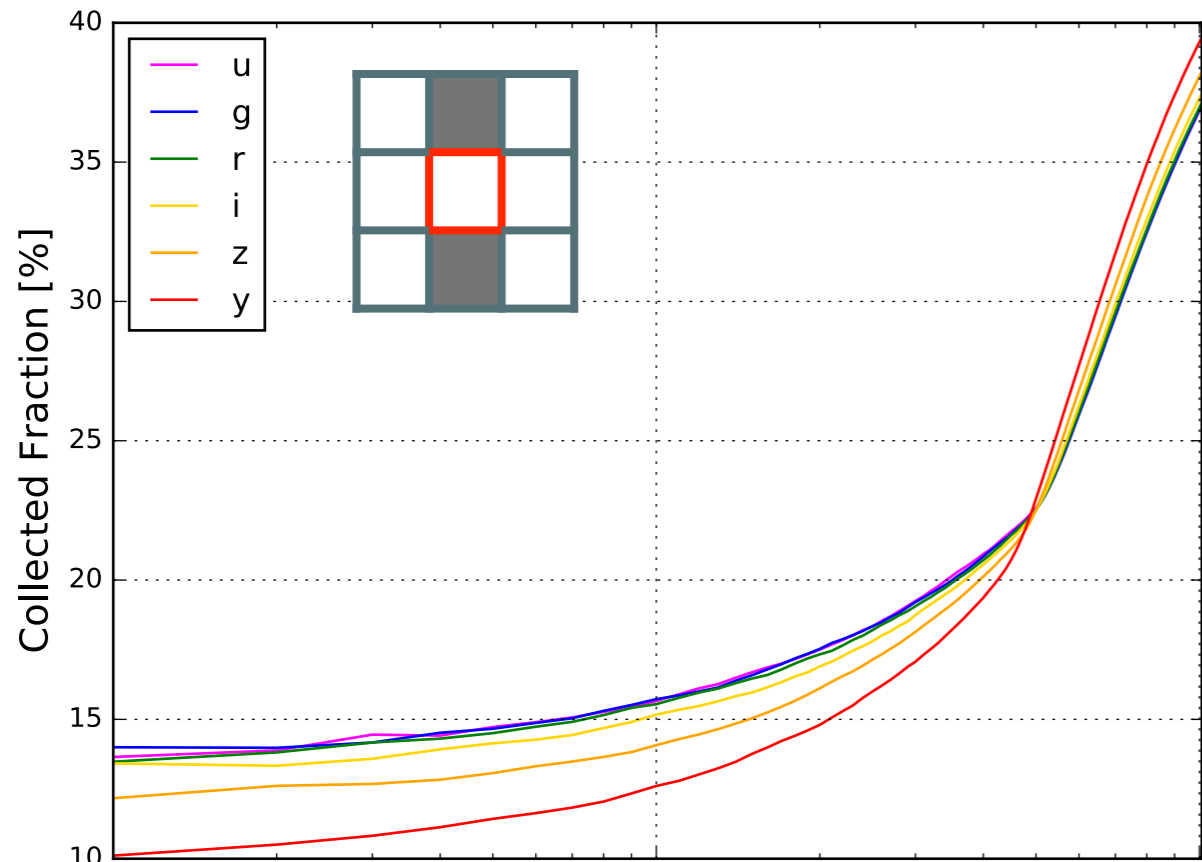












NEXT STEPS

- Compare with test bench data using tungsten lamp + filters and Fe-55 X-rays.
- Revisit how BF effect is modeled in GalSim?
 - Pixel effective area less well defined because of interplay of depth distribution with diffusion.
- Include effects of non-normal photon incidence and depth of focus.
 - Effects now depend on focal-plane position and break up/down & left/right symmetries.
 - Red sensor PSF is now bigger than blue!

BACKUP

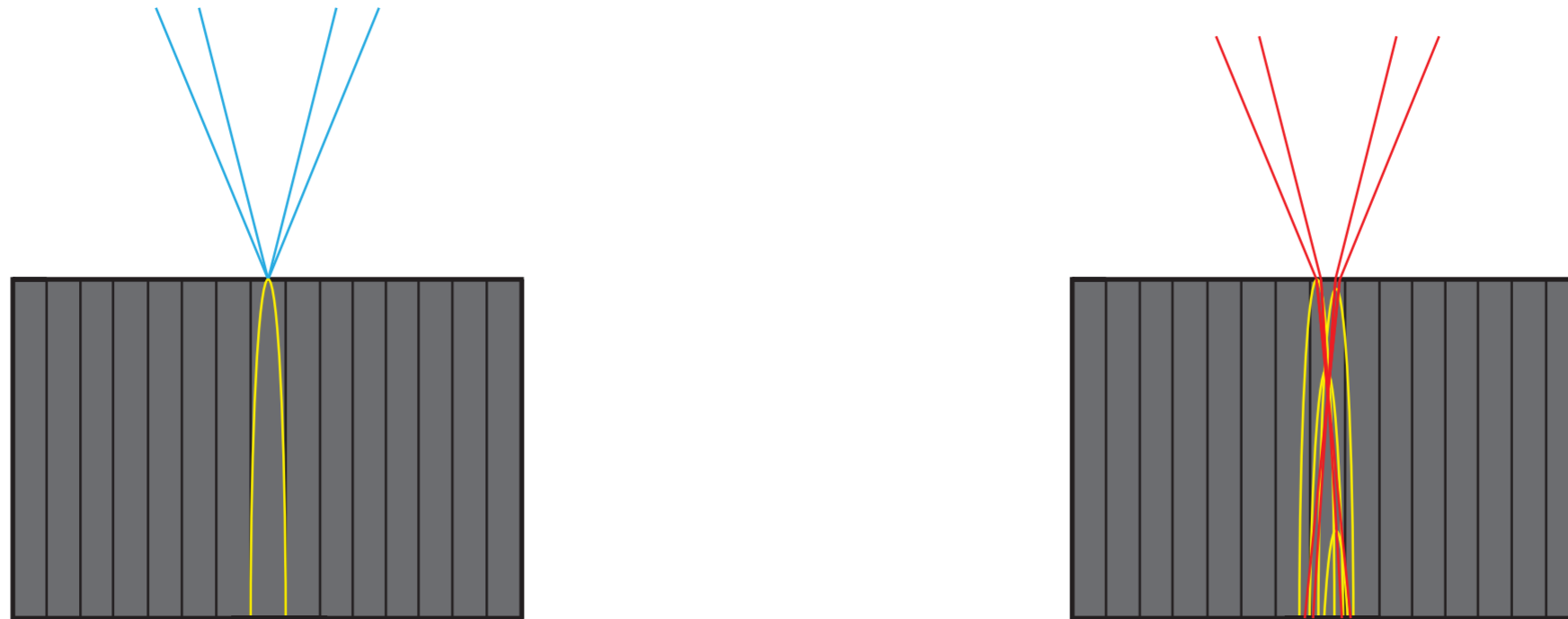


Figure 5. Cartoon figure of LSST beam and CCD geometry. Fifteen 10-micron wide and 100-micron deep pixels are shown in each cartoon. The $f/1.2$ beam is shown approaching the CCD from above for bluer filters (left) and redder filters (right). Blue photons convert into electron-hole pairs almost immediately, and the photo-electrons diffuse the entire depth of the CCD (shown in yellow). Red photons may penetrate deeper into the CCD bulk before converting. Combined with the diverging beam, which is still $f/4.5$ inside the silicon, the longer absorption length of redder photons leads to a larger PSF. Note that the position of best focus is a competition between beam divergence and charge diffusion. For bluer filters, the best focus is at the CCD surface, while for redder filters it is inside the bulk of the CCD [4].

Meyers & Burchat 2015 “Chromatic CCD effects on weak lensing measurements for LSST”, arxiv.org/abs/1505.02307

