### Observational Signatures of the Gammrays from Bright Blazars and Wakefield Theory

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# K. Abazajian, N. Canac, T. Tajima, T. Ebisuzaki, and S. Horiuchi, UC Irvine



## abstract

- **1.** Blazars emit γ-rays (observed: Fermi sat, etc.)
- 2. Blazars: AGN (blackhole) and their jets
- 3.  $\rightarrow$  time variations show: luminosity *L* and the power index *p* anti-correlate
- 4. Wakefield in the AGN jets triggered by disk disruptions: episodic accelerator of high energy electrons (thus  $\gamma$  photons) and UHECR
- 5.  $\rightarrow$  Wakefield acceleration: <u>anti-correlated</u> L and p

# Blazars in the Unified Model of AGN



### Core of Galaxy NGC 4261

Hubble Space Telescope

Wide Field / Planetary Camera

Ground-Based Optical/Radio Image HST Image of a Gas and Dust Disk 17 Arc Seconds 380 Arc Seconds **400 LIGHT-YEARS** 88,000 LIGHT-YEARS

## Anti-correlation between the luminosity and the power index from blazars



FIG. 2.— Shown are the flux (blue circles, left axis) and spectral index (green squares, right axis) for 3C 454.3 in 300 time bins of 7.9 days duration. An anti-correlation can be seen: the peaks in flux correspond to dips in the spectral index and vice versa.



Luminosity *L* and Power index *p* in time

Power index *p* vs. Luminosity *L* for several Blazars (more in Abazajian et al. arXiv)

# Luminosity of gamma ray emission and the spectrum AGN 3C454.3 with M = $10^7 M_{\odot}$



## Laser Wakefield (LWFA):

Wake phase velocity >> water movement speed <u>maintains</u> coherent and smooth structure



VS

Tsunami phase velocity becomes ~0, causes **wavebreak** and **turbulence** 



Strong beam (of laser / particles) drives plasma waves to saturation amplitude:  $E = m\omega v_{ph}/e$ No wave breaks and wake <u>peaks at v≈c</u> Wave breaks at v<c





**Relativistic coherence** enhances beyond the Tajima-Dawson field  $E = m\omega_p c/e$  (~ GeV/cm)

### Wakefield excited on the jets excited by MRI-driven disk disruption from BH: genesis of EHECR and gamma bursts



(28)

(29)

(31)

### Gamma emission luminosity by wakefield

 $L\gamma \sim 10^{33}$  ( $\kappa / 0.1$ )  $\eta$  m' m (erg/s)

 $\kappa$  (efficiency),  $\eta$  (episode dependent ~1)

#### (Ebisuzaki & Tajima: Astropart. Phy., 2014)

#### Fig. 2. (a the magn wave pul accelerat wakefield

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Fig. 3. P

strong pondermotive force. Eq. 25 holds as far as Zacc is greater than D. The distance D3 is where the acceleration finishes, defined by the equation  $D_3 = Z_{pd} = ac/\omega_A.$ 

We find that particles arrive at  $D_1$  before  $D_3$ , in other words:  $D_3/3R_g = 3.9 \times 10^5 (\dot{m}/0.1)^{5/3} (m/10^8)^{1/3} > D_1/3R_g.$ 

The energy spectrum of the accelerated charged particles has

the power-law with the index of -2 in the 1-D model due to the multiple dephasing occurrences when particles ride on and off different peaks of the pondermotive or wakefield hills when the waves contain multiple frequencies (but with again the same phase velocity ~ c; [8]), i.e.,  $f(W) = A(W/W_{min})^{-2}$ . As noted earlier, when the driving Alfven waves and their driven pondermotive fields hold a broad band of frequencies, their phase velocities and group velocities, respectively, are again close to the speed of light, providing the basis for the robust accelerating structure. When Alfven waves have two or three dimensional features, the dephasing is more prompt, leading to higher index of the spectrum (less than -2). Let κ be the energy conversion efficiency of the acceleration (including the mode convergence efficiency mentioned earlier), then  $\kappa E_B = AW_{\min}^2 ln(W_{\max}/W_{\min})$ , i.e. (30)

 $A = 1.6 \times 10^{33} \kappa \dot{m} m^2 [W_{\min}^2 ln(W_{\max}/W_{\min})]^{-1}$ 

The recurrence rate v<sub>A</sub> of the Alfven pulse burst is evaluated as:

 $v_A = \eta V_{AD}/Z_D = 1.0 \times 10^2 \eta m^{-1} \text{ Hz},$ where  $\eta$  is episode-dependent, and on the order of unity. This is

consistent with the 3-dimensional simulations conducted by O'Neill [12]. They found magnetic fluctuations, called Long Period Quasi Periodic Oscillations (LPQPO) with the period 10-20 times the Kepler rotation period. The luminosity L<sub>UHECR</sub> of ultra-high energy cosmic rays is:



Fig. 4. The total luminosities of accreting blackholes are plotted against the blackhole mass (in the unit of solar mass) for various maximum attainable energy  $W_{max}$  (solid lines) for the case of  $\Gamma = 20$  and  $\xi = 10^{-2}$ . Dashed lines are drawn for the values  $\dot{m} = 10^{-5}$ ,  $10^{-3}$ , and  $10^{-1}$ . The grey triangle represents the parameter set which allow the acceleration of UHECRs (  $\ge 10^{20}$  eV). We set the upper limit of  $\dot{m}$  to be around 0.1 for the pondermotive/wakefield acceleration to work, since the accretion disk becomes radiation dominant as *m* approaches unity, and the Alfven wave pulse becomes weaker than the estimate in the present pape

#### 4. Astrophysical implications and blazar characteristics

Radio galaxies belong to one category of AGN, which has radio lobes connected to the nucleus by relativistic jets. Their central engines are accreting supermassive ( $m = 10^6 - 10^{10}$ ) blackholes. Urry and Padovani [27] pointed out that there are parent (or misaligned)



# Time evolution log10(plasma beta= $P_{Gas}/P_{Mag}$ ) @ equatorial plane





Plasma beta ~1
near the horizon.
(short time variability =~10)
anti-symmetric mode
stratified structure

# Summary

- Observe → Anti-correlation of time series of Blazars: Luminosity
   L(t) vs. Power Index p(t)
  - For Blazar whose mass is known (or guessed) the above rise and interval times proportional to the central mass, whose <u>absolute times</u>: in reasonable agreement below
- Theory → Wakefield acceleration theory predicted the same above Unlike Fermi mechanism difficulties, <u>no limit</u> in energies (no synchrotron, no confinement limits)
- 3. General Relativistic MHD simulation → episodic disruption whose time scales in agreement with the observed time series features above

# Thank you!